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**Design of ultra-sensitive fluorescence detection system for applications in molecular
biology.**

Thesis Presented

by

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to

The Graduate School

in Partial Fulfillment of the

Requirements

for the Degree of

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Abstract of the Thesis

Design of ultra-sensitive fluorescence detection system for applications in molecular biology.

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2012

I present design of a new ultra-sensitive data acquisition system for fluorescence detection applications. Based on single photon detection, the developed system comprises a unique single photon spectrometer with high detection speed and wide linearity dynamic range, and a fiberized capillary sample analyzer.

The system is capable of detecting multicolor fluorescence produced by mixtures of fluorescent dyes as well as by micro and nano-objects. In the system, the analyzed sample is forced into the capillary by applying either air pressure or electric field. When the sample passes through a detection area, a laser beam induces multi-color fluorescence. Polychromatic fluorescence is collected by a fiber and directed to the single photon spectrometer. The spectrometer performs color separation and the measurement of fluorescence in the range of wavelengths from 480nm to 720nm. The fluorescence which comes to the spectrometer through the input fiber is coupled to a collimator which produces a parallel polychromatic beam of 10mm diameter. The parallel beam passes through laser rejection filters and undergoes separation on the diffraction grating into constituent wavelength components. The separated monochromatic beams are focused onto channels of the 32-channel PMT by the system comprising spherical mirror. The received photons produce very short current pulses in channels of the 32-channel PMT. The pulses undergo amplification and counting. The obtained photon count is transferred to computer for recording and data processing. The unique combination of fluorescent dyes with

known spectra is recognized in each individual measurement by using color deconvolution of the dye's spectra.

The developed data acquisition system and its component modules can be used for many applications in Life Sciences, Physics, and other fields.

I applied the developed data acquisition system for detection of multicolor fluorescence of mixtures of quantum dots, high speed detection of micro-beads encoded with quantum dots and fluorescent dyes, and DNA sequencing by capillary electrophoresis.

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List of Abbreviations

Abbreviation or Symbol	Term
APD	Avalanche Photodiode
CE	Capillary Electrophoresis
DC	Dark Counts
OD	Optical Density
LIF	Laser Induced Fluorescence
PCB	Printed Circuit Board
PCR	Polymerase Chain Reaction
PMT	Photomultiplier Tube
SiPM	Silicon Photomultiplier
SPAD	Single Photon Avalanche Diode

I. INTRODUCTION

One of very important practical tasks in molecular biology is creation of multiplex measurements, when multiple biological reporters can be processed and measured within one experiment. In order to introduce multiplexing, it is common to use multi-color fluorescent markers for biological molecules. Another very efficient method for multiplex measurements is the use of color-encoded micro-beads. In this case the color encoding can be introduced due to the use of various compositions of several fluorescence dyes. For example, company Luminex offers up to 500 various types of encoded beads which are colored with only three colors. Recently a number of technologies have been developed for optical “bar coding” of micro- and nano-objects (beads, cells, etc.) using multi-color fluorescence. This includes the use of segmented nanorods [Nicewarner-Pena, S. R, 2001], rare-earth doped glass [Dejneka, M. J, 2003], fluorescent silica colloids [Grondahl, L, 2000], photobleached patterns [Braeckmans, K., 2003], enhanced Raman nanoparticles [Mulvaney, S. P., 2003] and semiconductor quantum dots (QD) [Gao, X.; Nie, S., 2003, 2004].

The goal of my work is design and development of a novel data acquisition system which allows a highly sensitive, accurate and fast detection of multicolor, laser induced fluorescence (LIF). The applications of our system will be in measurements and characterization of various biological molecules carried in solutions by different means (electrophoresis, fluidic pumping, molecules carried on micro-beads, etc.). Biological applications of our system will be detection and quantification of bio-molecules (proteins, RNA, DNA), which is very important in molecular biology and medicine.

The main goal of my work

My main goal was to design and develop a system which would be able to detect thousands of color combinations due to the use of single photon spectrometer in conjunction with appropriate fluidic system and software.

Detailed description of the system modules is presented in Chapters 1-2, and in Chapter 3 describes applications of the developed system for measurement of composition of fluorescent dyes and measurements of beads encoded with fluorescent dyes and quantum dots.

Operation of the system

The system consists of two main modules: single photon spectrometer and capillary sample analyzer, and an external laser source (in my work I used various commercial lasers which were fiberized with either single-mode or multi-mode fibers). The architecture of the system is such that the modules are connected via optical fibers and the detected signals are transferred to computer for recording and processing (see Figure I.1). In the system a fiberized laser source delivers excitation power to a capillary sample analyzer. In the analyzer, fluorescent objects are moved through the capillary in which they are illuminated by the laser, their fluorescence is collected by a fiberized optical collection system and delivered to a fiberized spectrometer. In the spectrometer, after color decomposition the fluorescent signal is detected by a multi-channel single photon detector, the detected signals are transferred to the computer.

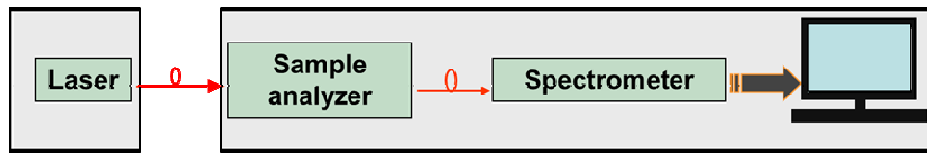


Figure I.1. Architecture of the system

Approach to system design

Our approach to system design is presented in Figure.I2.

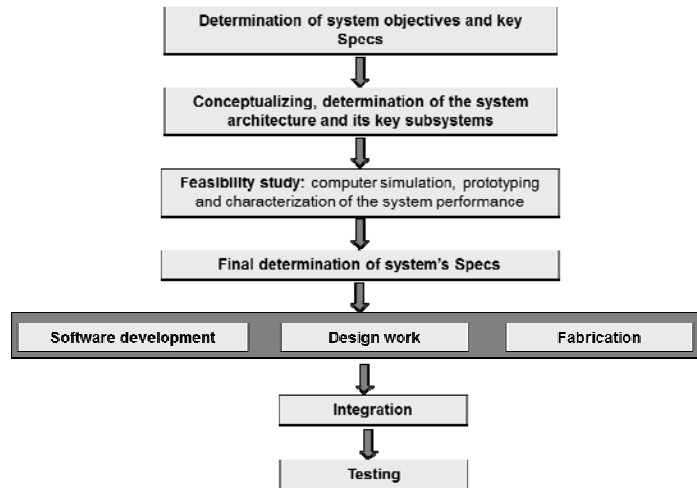


Figure I2. Approach to system design and development.

1. CHAPTER 1. SPECTROMETR

1.1.System objectives and key specs

Our system is intended for highly sensitive, accurate and fast detection and recognition of fluorescence encoded micro- and nano-objects. Since our system is intended to detect very weak fluorescent signals, we select single photon counting as a detection method.

In order to have a high detection throughput, we need to be able to detect up to 10,000 events per second (~20mks/event). In order to ensure high detection accuracy, our spectrometer must have broad linearity range and high frame rate.

The target specs for the spectrometer are presented below:

Spectral range	520-750nm (<i>determined by fluorescent dyes used</i>)
Spectral resolution	5-10nm/channel (<i>based on spectral width of fluorescent dyes ~50nm</i>)
Linearity:	10 ⁸ /s per channel
Frame rate:	1MHz
Data recording rate:	32MB/s via USB port

Operation of the spectrometer

Block-diagram of single photon spectrometer is presented in Figure 1.1.1.

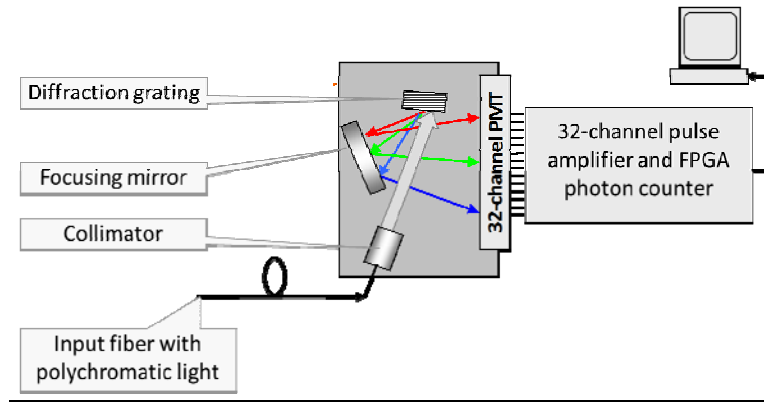


Figure 1.1.1. Schematic of spectrometer

In the spectrometer polychromatic light delivered through optical fiber, passes through collimator, gets decomposed by a diffraction grating, and focused onto photocathodes of 32-channel PMT. The signals detected by the PMT are amplified by high frequency pulse amplifier, counted by FPGA photon counter and transferred to computer.

1.2.Theoretical background

Calculations

Diffraction grating

The diffraction grating is of considerable importance in spectroscopy, due to its ability to separate (disperse) polychromatic light into its constituent monochromatic components. In recent years, the spectroscopic quality of diffraction gratings has greatly improved, and Newport has been a leader in this development.

The extremely high accuracy required of a modern diffraction grating dictates that the mechanical dimensions of diamond tools, ruling engines, and optical recording hardware, as well as their environmental conditions, be controlled to the very limit of that which is physically possible. A lower degree of accuracy results in gratings that are ornamental but have little

technical or scientific value. The challenge to produce precision diffraction gratings has attracted the attention of some of the world's most capable scientists and technicians. Only a few have met with any appreciable degree of success, each limited by the technology available.

A *diffraction grating* is a collection of reflecting (or transmitting) elements separated by a distance comparable to the wavelength of light under study. It may be thought of as a collection of diffracting elements, such as a pattern of transparent slits (or apertures) in an opaque screen, or a collection of reflecting grooves on a substrate (also called a *blank*). In either case, the fundamental physical characteristic of a diffraction grating is the spatial modulation of the refractive index. Upon diffraction, an electromagnetic wave incident on a grating will have its electric field amplitude, or phase, or both, modified in a predictable manner, due to the periodic variation in refractive index in the region near the surface of the grating.

A *reflection grating* consists of a grating superimposed on a reflective surface, whereas a *transmission grating* consists of a grating superimposed on a transparent surface.

A *master grating* (also called an *original*) is a grating whose surface-relief pattern is created “from scratch”, either by mechanical ruling or holographic recording. A *replica grating* is one whose surface relief pattern is generated by casting or molding the relief pattern of another grating.

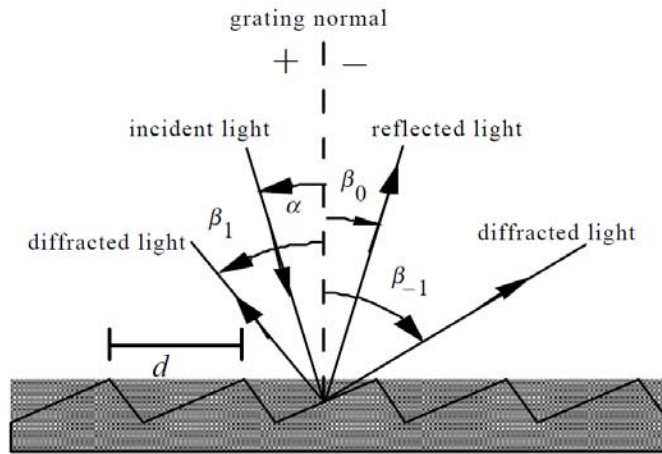


Figure 1.2.1. Diffraction grating ray geometry.

When a grating depends on the fact that there exists a unique set of discrete angles along which, for a given spacing d between grooves, the diffracted light from each facet is in phase with the light diffracted from any other facet, leading to constructive interference.

Diffraction by a grating can be visualized from the geometry in Figure 1.2.1 , which shows a light ray of wavelength l incident at an angle a and diffracted by a grating (of *groove spacing* d , also called the *pitch*) along at set of angles $\{bm\}$. These angles are measured from the grating normal, which is shown as the dashed line perpendicular to the grating surface at its center. The sign convention for these angles depends on whether the light is diffracted on the same side or the opposite side of the grating as the incident light. In Figure 1.2.1 which shows a *reflection grating*, the angles $a > 0$ and $b_1 > 0$ (since they are measured counter-clockwise from the grating normal) while the angles $b_0 < 0$ and $b_{-1} < 0$ (since they are measured clockwise from the grating normal)

Spectrometer calculations

Schematic of the optical system of the spectrometer is shown in Figure 1.2.2.

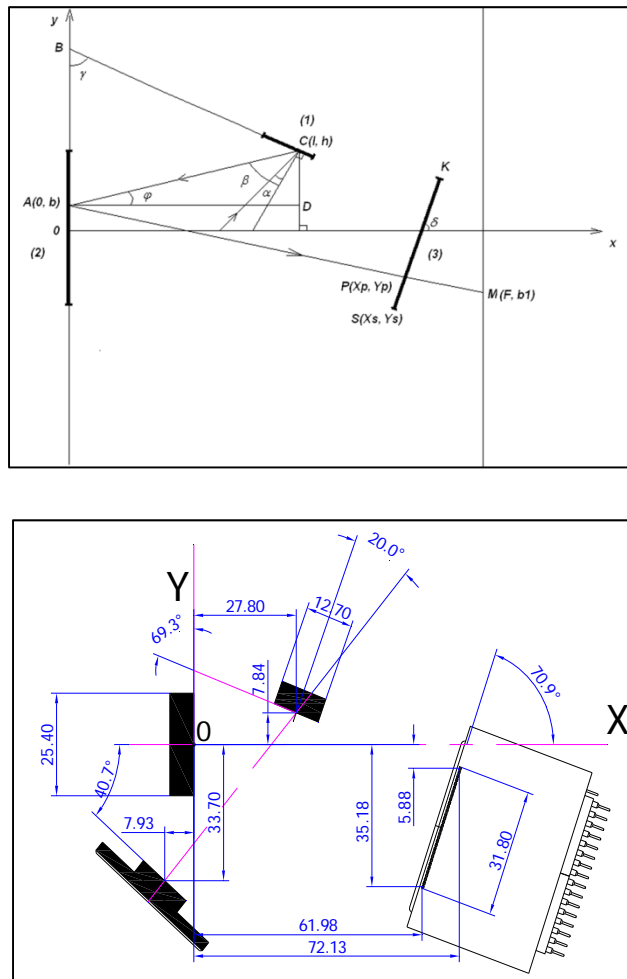


Figure 1.2.2. Schematic and measurements of the spectrometer.

Polychromatic light impinged on the grating (1) has incident angle α , reflected on mirror (2) has angle β , reflected one more time on detector and has angle F .

Positions and dimensions of diffraction grating, mirror and detector are given in Figure 1.2.2 (upper panel).

Now we need to find which channel of the detector is illuminated by the back light ray.

1. Let's find correlation between angles α, β, φ .

We know equation for angles correlation between incident light and reflected light on diffraction grating

$$\beta = \arcsin\left(\frac{m\lambda}{d_{DG}} - \sin \alpha\right), \quad (1)$$

Where

$$m = 0, \pm 1, \pm 2, \dots,$$

α - wavelength

d_{DG} - parameters of diffraction grating.

From the drawing we can see that:

$$\varphi + \beta + \frac{\pi}{2} + \pi - \gamma = \pi$$

Therefore,

$$\varphi = \gamma - \beta - \frac{\pi}{2}.$$

From triangle $\square ACD$ we can find b :

$$\operatorname{tg} \varphi = \frac{h-b}{l}, \text{ from her we can find next equation}$$

$$b_{\beta} = h - l \cdot \operatorname{tg} \varphi . \quad (2)$$

We know that $-\operatorname{tg} \varphi = \frac{b_1}{F}$, then

$$b_1 = -F \cdot \operatorname{tg} \varphi . \quad (3)$$

2. Let's find point where incident light beam crosses axle OY .

Angle between incident ray and positive direction OX , equal $\gamma - \frac{\pi}{2} - \alpha$, and crosses point $C(l, h)$ on grating. Then equation of incident ray will be:

$$y = -x \cdot \operatorname{ctg}(\gamma - \alpha) + h + l \cdot \operatorname{ctg}(\gamma - \alpha) .$$

Let's call coordinate point where incident light ray and axe OY as point b_{α} .

Then we have

$$b_{\alpha} = h + l \cdot \operatorname{ctg}(\gamma - \alpha)$$

3. Now we'll find cross point reflected light ray from mirror and detector

This ray crosses two point that we know $A(0, b)$ and $M(F, b_1)$.

From here we have:

$$y = \frac{b_1 - b}{F} x + b .$$

Detector is on the line that creates angle δ , with axe OX and crosses point $S(x_s, y_s)$.

Then equation for the line that crosses detector:

$$y = x \cdot \operatorname{tg} \delta + Y_s - X_s \cdot \operatorname{tg} \delta$$

Now we need to find cross point between these two lines. To do that we need to solve equation:

$$\begin{cases} y = \frac{b_1 - b}{F} x + b \\ y = x \cdot \operatorname{tg} \delta + Y_s - X_s \cdot \operatorname{tg} \delta \end{cases}$$

Solution is:

$$X_{PMT} = \frac{X_s \cdot \operatorname{tg} \delta - Y_s + b}{F \cdot \operatorname{tg} \delta - b_1 + b} \cdot F \quad (4)$$

$$Y_{PMT} = \frac{(b_1 - b) \cdot (X_s \cdot \operatorname{tg} \delta - Y_s + b)}{F \cdot \operatorname{tg} \delta - b_1 + b} + b$$

Now we find $PS = \sqrt{(X_s - X_{PMT})^2 + (Y_s - Y_{PMT})^2}$ (5).

Point K belongs to the first channel and point S - to the channel number 32.

$$\frac{32(q-PS)}{q}$$

Then point P belongs $\frac{32(q-PS)}{q}$ channel, where q - detector length. (6)

4. Now we need to find cross point of the reflected ray from the mirror, and diffraction grating.

Diffraction grating is located on a line that has creates angle γ with axe OX, and crosses through the middle of diffraction grating that has coordinates (l_0, h_0)

Now we can find equation for the line that crosses diffraction grating:

$$y = x \cdot \operatorname{tg} \gamma + h_0 - l_0 \cdot \operatorname{tg} \gamma$$

In order to find cross point between these two lines we need to solve these equations:

$$\begin{cases} y = \frac{b_1 - b}{F} x + b \\ y = x \cdot \operatorname{tg} \gamma + h_0 - l_0 \cdot \operatorname{tg} \gamma \end{cases}$$

Solution is:

$$x_{DG} = \frac{l_0 \cdot \operatorname{tg} \gamma + h_0 - b}{F \cdot \operatorname{tg} \gamma + b_1 - b} \cdot F \quad (7)$$

$$y_{DG} = \frac{(b_1 - b)(l_0 \cdot \operatorname{tg} \gamma + h_0 - b)}{F \cdot \operatorname{tg} \gamma + b_1 - b} + b$$

These points should satisfy next:

$$\begin{cases} l_{\min} \leq X_{DG} \leq l_{\max} \\ h_{\min} \leq Y_{DG} \leq h_{\max} \end{cases} \quad (8),$$

Where $l_{\min}, h_{\min}, l_{\max}, h_{\max}$ determined as.

$$C_{\max} : \begin{cases} l_{\max} = l_0 - \left(-\frac{D}{2} \cos \gamma \right) = l_0 + \frac{D}{2} \cos \gamma \\ h_{\max} = h_0 + \frac{D}{2} \sin \gamma = l_0 + \frac{D}{2} \sin \gamma \end{cases} \quad (9)$$

$$C_{\min} : \begin{cases} l_{\min} = l_0 + \left(-\frac{D}{2} \cos \gamma \right) = l_0 - \frac{D}{2} \cos \gamma \\ h_{\min} = h_0 - \frac{D}{2} \sin \gamma = l_0 - \frac{D}{2} \sin \gamma \end{cases}$$

If reflected ray crosses diffraction grating, then we can use next equations:

$$\left\{ \begin{array}{l} x_{DG} = l_{\min} \\ \frac{b-b_1}{F} l_{\min} + b - h_{\min} \geq 0 \end{array} \right. \quad \text{and} \quad \left\{ \begin{array}{l} x_{DG} = l_{\max} \\ \frac{b-b_1}{F} l_{\max} + b - h_{\max} \leq 0 \end{array} \right. \quad (10)$$

5. Now let's say we have the light ray, with specific wavelength, that crosses point C_{\min} with angle α , and angle $\varphi = 0$.

Now we need to determine angle α that satisfies $\varphi = 0$.

We know

$$\varphi = \gamma - \beta - \frac{\pi}{2},$$

than $0 = \gamma - \beta - \frac{\pi}{2} \Rightarrow \beta = \gamma - \frac{\pi}{2}$.

Therefore,

$$\beta_{\varphi=0} = \gamma - \frac{\pi}{2} = \arcsin\left(\frac{m\lambda}{d_{DG}} - \sin \alpha_{\varphi=0}\right), \quad \left| \frac{m\lambda}{d_{DG}} - \sin \alpha_{\varphi=0} \right| \leq 1$$

We need to solve this equation:

$$\sin\left(\gamma - \frac{\pi}{2}\right) = \sin\left(\arcsin\left(\frac{m\lambda}{d_{DG}} - \sin \alpha_{\varphi=0}\right)\right)$$

$$-\cos \gamma = \frac{m\lambda}{d_{DG}} - \sin \alpha_{\varphi=0}$$

$$\sin \alpha_{\varphi=0} = \frac{m\lambda}{d_{DG}} + \cos \gamma$$

$$\alpha_{\varphi=0} = \arcsin \left(\frac{m\lambda}{d_{DG}} + \cos \gamma \right), \quad \left| \frac{m\lambda}{d_{DG}} + \cos \gamma \right| \leq 1 \quad (11)$$

If we'll increase angle α , then some specific ray with specific wavelength will cross diffraction grating. For each wavelength we'll have different angle $\alpha_{\varphi=0}$. To avoid situation when reflected ray crosses diffraction grating, incident light ray should have angle that satisfy:

$$< \min \left\{ \alpha_{\varphi=0, \lambda_1}, \alpha_{\varphi=0, \lambda_2}, \alpha_{\varphi=0, \lambda_3}, \dots \right\}.$$

1.3.Design of 32-channel spectrometer

Based on the obtained equations we carried out preliminary calculations of the spectrometer.

Final calculations were carried our using CAD package (Figure 1.3.1)

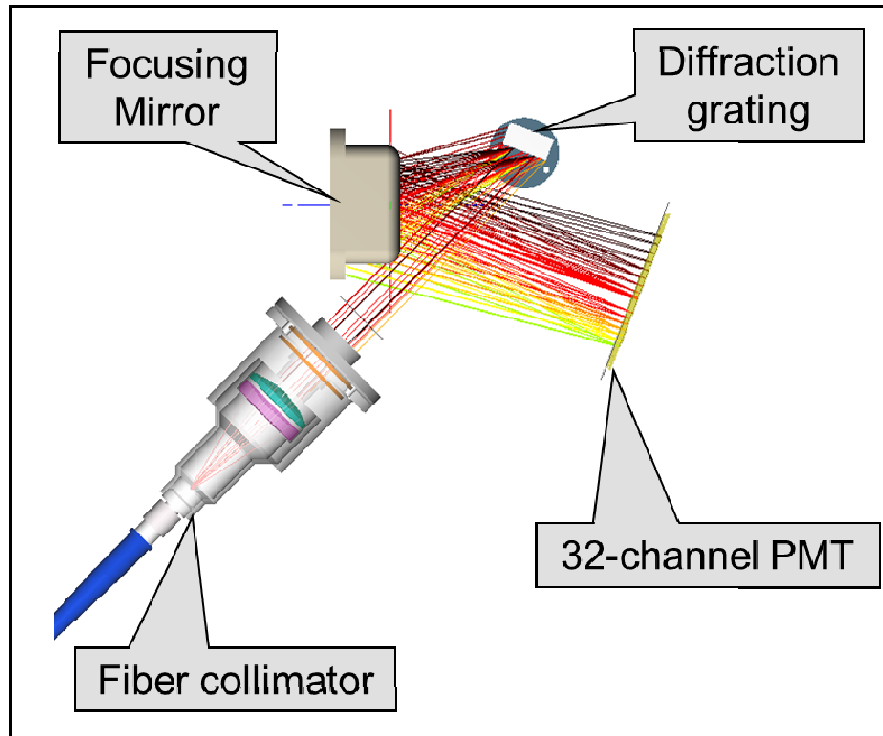


Figure 1.3.1. Ray tracing in spectrometer

The spectrometer was designed based on the following commercial components:

1. The SMA Optical fiber core diameter 200mkm
2. The F810SMA Series of fiber collimation
3. Thorlabs GR13-1850 - ruled diffraction grating, 1800/mm, 500 nm Blaze, 12.7 x 12.7 x 6 mm
4. CM254-050-G01 - Ø1" Aluminum-Coated Concave Mirror, f=50 mm
5. Filters

Collimator assembly

The F810SMA Series of fiber collimation packages is pre-aligned to collimate a laser beam propagating from the tip of an SMA-connectorized fiber with diffraction-limited performance at the design wavelength. Since the F810 Series fiber collimators do not have any movable parts, they are compact and not susceptible to misalignment. Due to chromatic aberration, the effective focal length (EFL) of the doublet lens is wavelength dependent. As a result, these collimators will only perform optimally at the design wavelength. The doublet lens is factory aligned so that it is one focal length away from the fiber tip when inserted into the collimator. This distance is equal to the focal length of the aspheric lens at the design wavelength. In addition, the doublet lens has an AR coating that minimizes surface reflections.

Ruled diffraction grating

These replicated, ruled diffraction gratings are offered in a variety of sizes and blaze angles. Ruled gratings typically achieve higher efficiencies than holographic gratings due to their blaze angles. Efficiency curves for all of these gratings are shown on the Figure 1.3.2

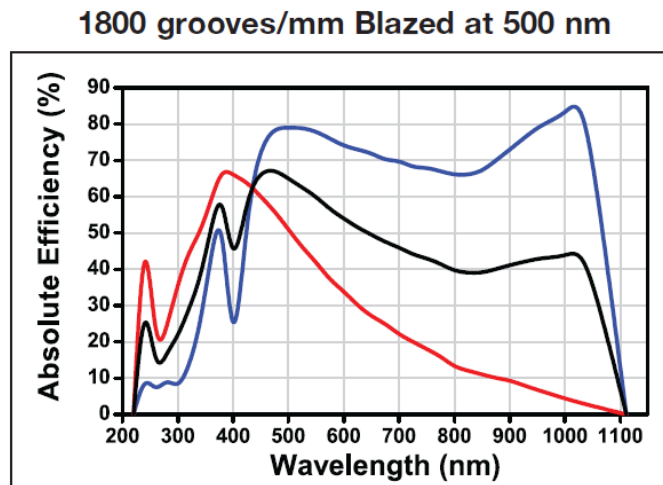


Figure 1.3.2 Absolute Efficiency

Concave mirror

Concave Mirrors offer high reflectance in broadband applications where the chromatic aberration introduced by a lens is not acceptable. When used close to 0° angle of incidence to collect and focus light for imaging applications, these mirrors provide $>90\%$ reflectance with the metallic coatings and $>99\%$ reflectance with the dielectric coatings. Spherical Concave Mirrors are defined by the focal length, f , of the optic being half of its radius of curvature, R . To achieve the best performance without the introduction of off axis aberrations, concave mirrors should be used close to 0° angle of incidence.

Specifications

- Focal Lengths: 12 mm – 500 mm
- Substrate Material: N-BK7
- Reflective Coatings: Metallic and Dielectric Coatings
- Clear Aperture: $>90\%$ of Diameter
- Front Surface Irregularity: $<\lambda/4$ at 633 nm
- Front Surface Quality: 40-20 Scratch-Dig
- Back Surface: Fine Ground
- Diameter Tolerance: $+0.0/-0.2$ mm
- Thickness Tolerance: ± 0.2 mm
- Chamfer: 0.25 mm @ 45° Typical

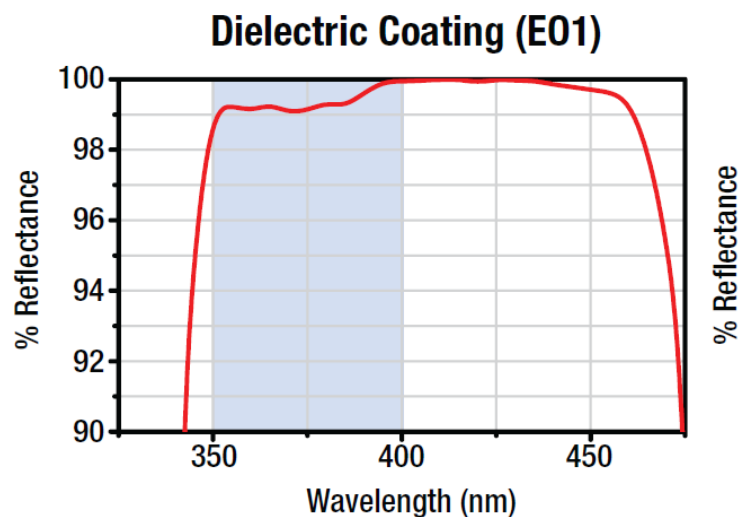


Figure 1.3.3. Reflectance of diffraction grating versus wavelength.

In Figure 1.3.3 the shaded region represents the specified region for each coating. However when used at near 0° AOI, as is recommended for a concave mirror, the wavelength range where the mirror has a high reflectance is significantly larger as can be seen in the performance

Optical filters

Longpass filters offer unprecedented performance in filter applications. The steep edges (measured from an optical density of 6.0 to a transmission of 50%) make it possible to measure even the smallest Raman shifts.

The large bandwidths and exceptional transmission permit these filters to be used in even the most demanding imaging applications.

One benefit of these laser filters, the rejected laser source light is reflected, easing stray-light management issues associated with holographic notch filters. Eliminating the need for complicated light baffling simplifies system design, creating more compact, cost-effective instrumentation. Second enhancing feature of these filters, with guaranteed transition widths (measured from the laser wavelength to the 50% transmission wavelength) of <1%, angle-tuning is unnecessary, greatly simplifying system setup and implementation. And finally, the coating process ensures unrivaled laser filter lifetime and durability, and guarantees “near-zero” temperature dependence, ensuring low long-term cost of ownership and maximum system operational range and sensitivity. These laser filters are not intended for general purpose longpass applications, as they do not have consistent blocking in the shortwave region.

Specifications

- Attenuation OD: 6
- Clear Aperture (mm): 21.3
- Filter Type: Longpass Filters
- Attenuation Range (NM): UV - Cut-on
- Exact Wavelength: 520
- Description: 520LP
- Brand: 3rd Millennium

- Peak T%: ≥ 90
- Cut-on (nm): 520
- Thickness (MM): 5.5

Figure 1.3.4 shows CAD of optical bench and assembly with PMT enclosure (upper panels) and assembly of spectrometer bench, PMT and PCB (lower panel).

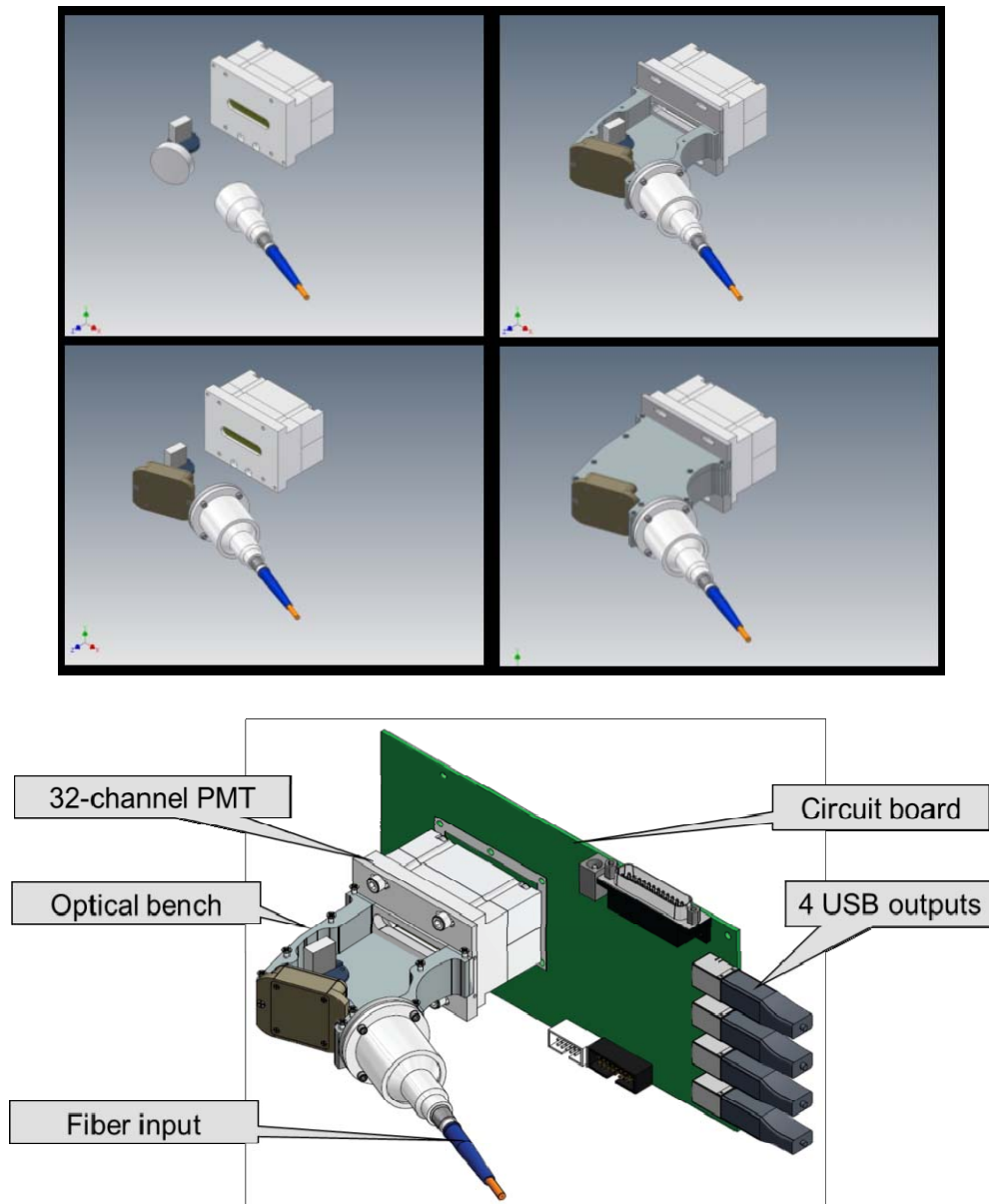


Figure 1.3.4. CAD of optical bench

As can be seen, the optical bench carries diffraction grating mounted on a fixture which has adjustable angle, Spherical mirror in a housing, fiberized collimator, and PMT housing with adjustable position. Lower panel of Figure 1.3.4 shows how optical bench is assemble with PCB. The PCB contains channel pulse amplifiers, 32-channel USB photon counter and has 4 USB outputs.

Figure 1.3.5 shows block diagram and CAD of the spectrometer assembly.

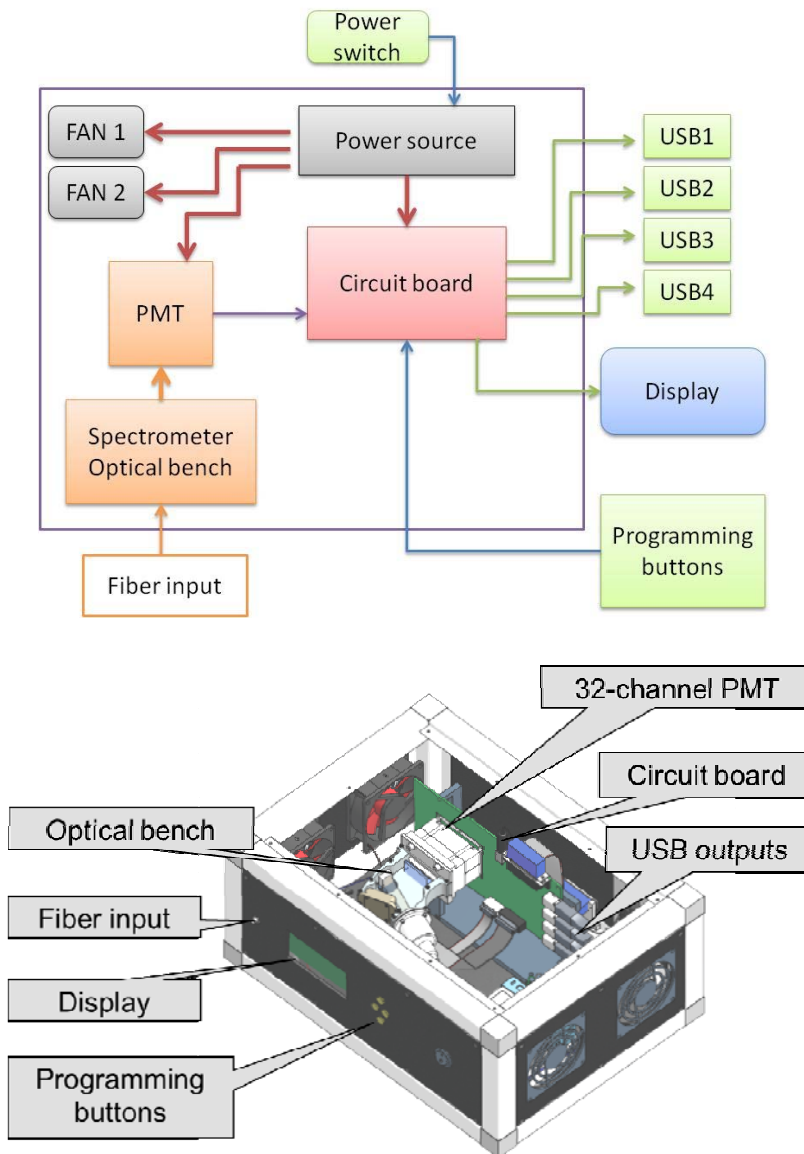


Figure 1.3.5. Block diagram and CAD of spectrometer assembly

The spectrometer assembly comprises a stabilized dc power supply (NES-100-5), which has voltage output 5V and maximum current 20A. Same voltage output supports two fans, PMT and used for supporting the PCB. The device is controlled by programming buttons which allow choosing different detection and recording regimes. The display shows the chosen operation regime. Optical signals are delivered to the spectrometer via input 200 μ m fiber.

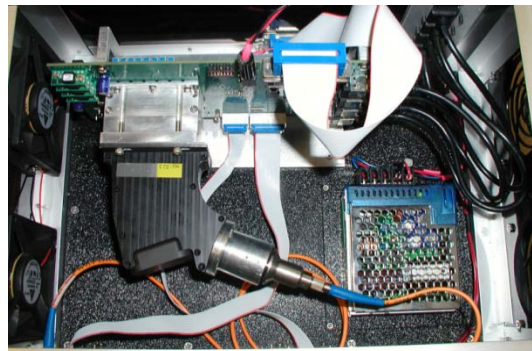
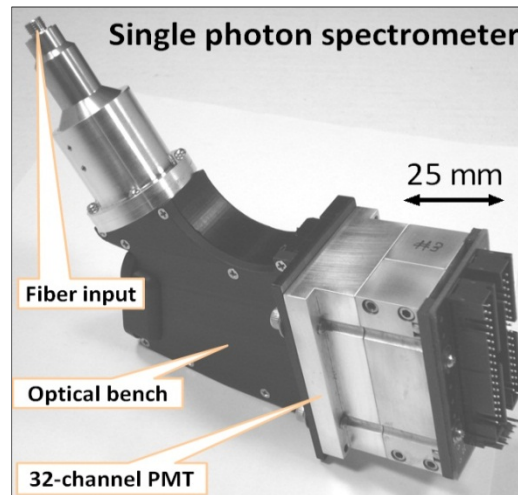


Figure 1.3.6. Photographs of the spectrometer

Figure 1.3.6 shows the finalized spectrometer. Upper panel presents a photograph of the optical bench assembly, inner view of the spectrometer assembly is shown on middle panel, and general view of the spectrometer is presented on lower panel.

1.4. Testing of 32-channel single photon spectrometer

In order to characterize the spectrometer we connected the input fiber of the spectrometer to a monochromator (Figure 1.4.1) and carried out measurements of the spectrometer response to different wavelengths.

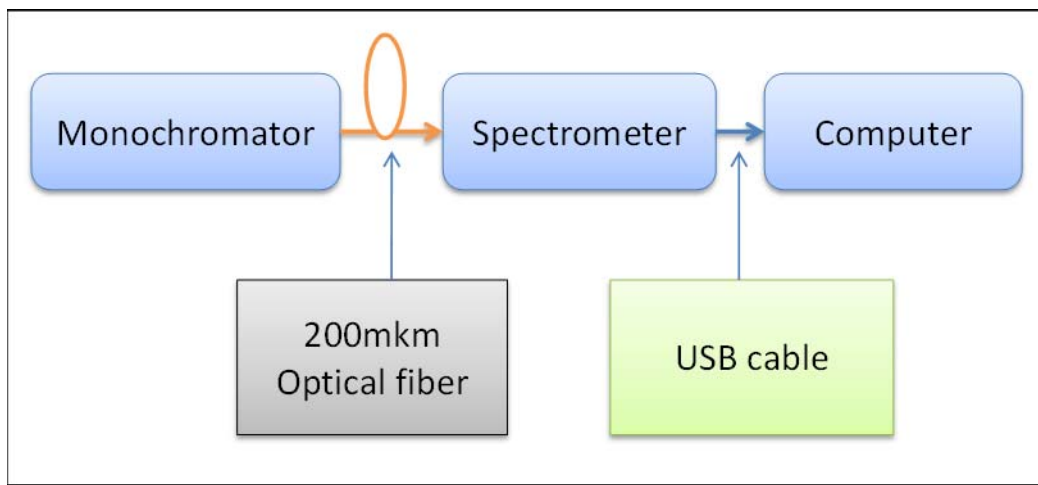


Figure 1.4.1. Block-diagram of the spectrometer characterization setup

Figure 1.4.2 presents channel resolution (upper panel) and wavelength resolution (lower panel) of the spectrometer with 32-channel PMT sensor. As can be seen, wavelengths which differ by 10nm are mostly detected by one channel of the PMT. Lower panel of the Figure 1.4.2 shows the wavelength resolution of the sensor. As can be seen, each PMT channel detects wavelengths in the range of $\sim 10\text{nm}$.

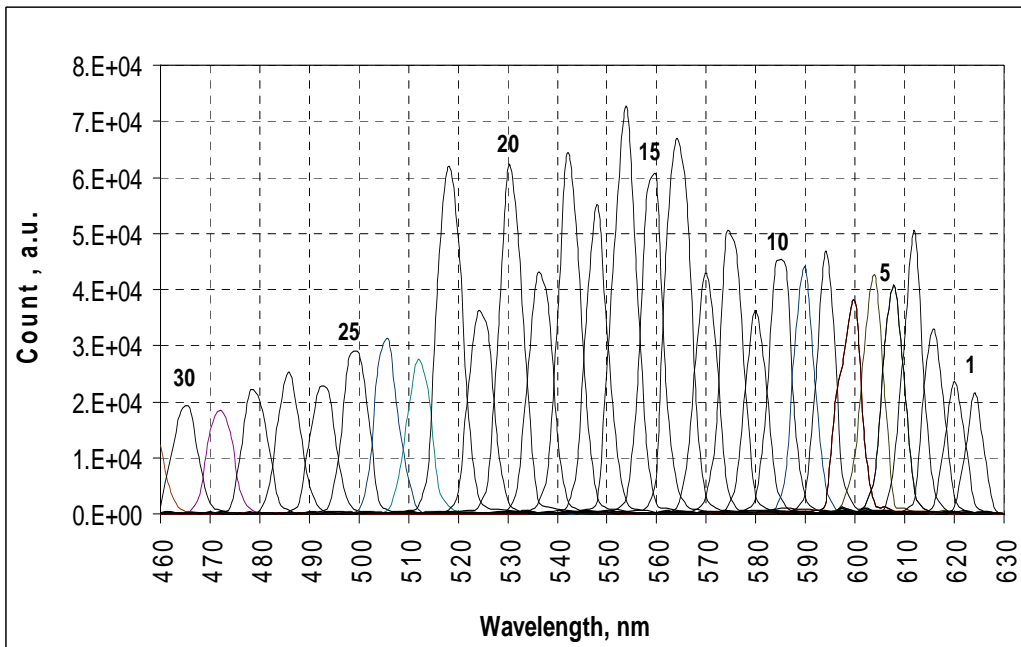
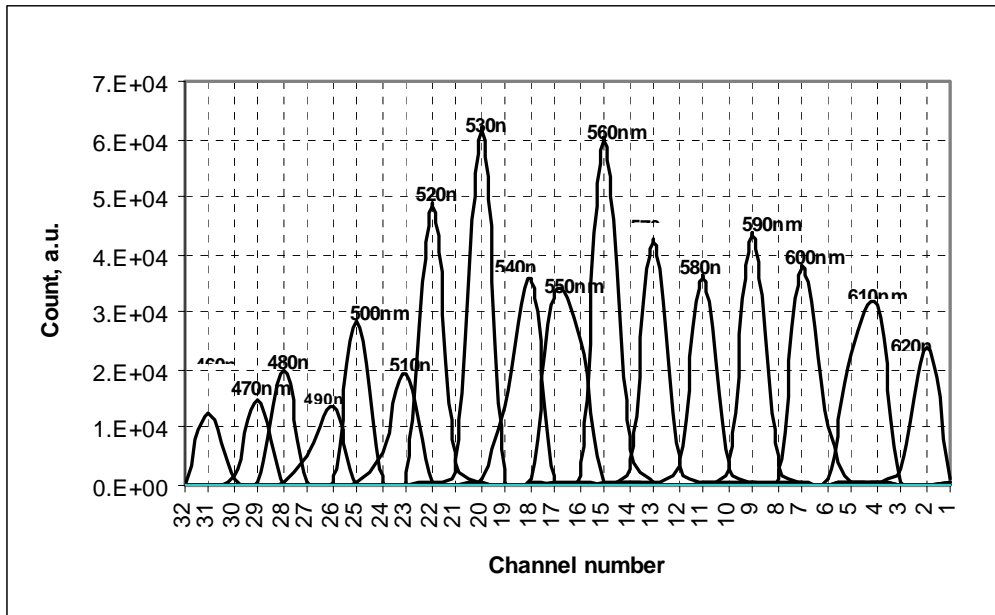


Figure 1.4.2. Spectral resolution of the spectrometer

Analysis of the spectral separation module of the spectrometer shows that it can provide spectral resolution as high as 1nm. The obtained spectral resolution of $\sim 10\text{nm}$ is determined by the geometry of the 32-channel PMT which comprises $0.8\text{mm} \times 7\text{mm}$ detection zones separated by 0.2mm distance.

The overall resolution of the sensor can be improved by using an array of photoreceiving fibers connected to single photon sensors (see photos in Figure 1.4.3).

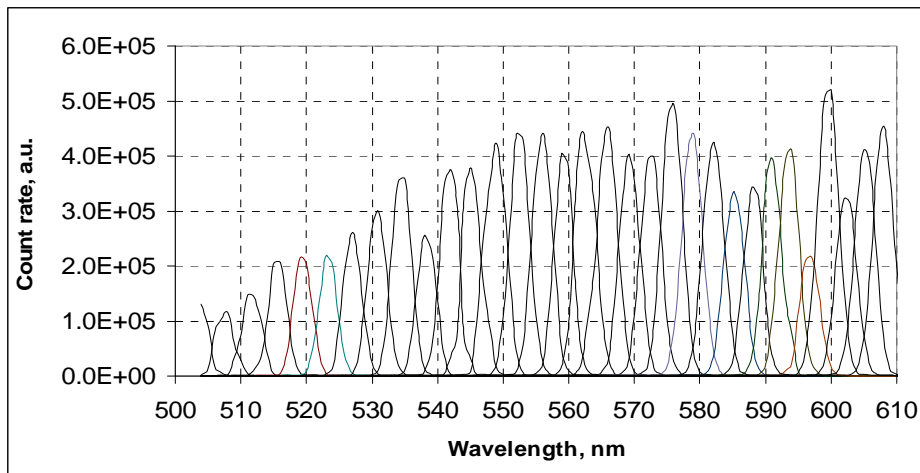
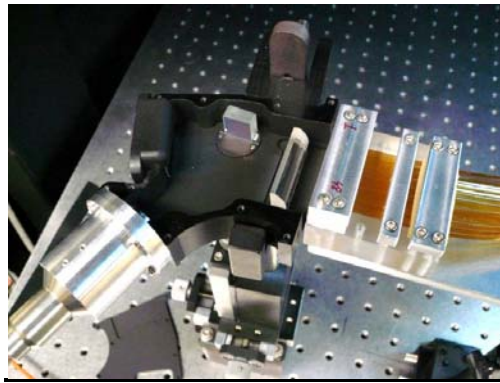
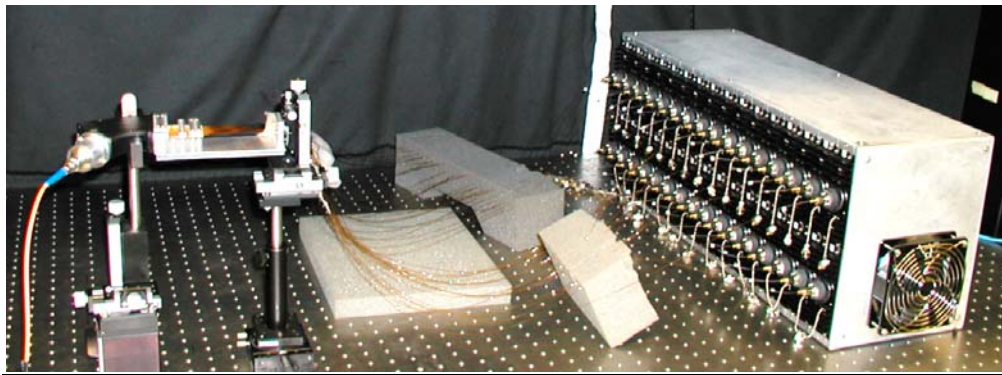


Figure 1.4.3. Spectrometer with array of photo-receiving fibers (400µm core, 650µm pitch, upper and middle panels) and wavelength resolution (lower panel).

As can be seen on the lower panel of Figure 1.4.3, the spectral resolution of the spectrometer with the fiber bundle is ~ 5 nm. In fact, it means that the receiving fibers can be used as band pass filters in some applications when high spectral resolution is required.

2. CHAPTER 2. CAPILLARY FLUORESCENCE ANALYZER

2.1. System objective and key specs

The main goal of the capillary fluorescence analyzer is to deliver fluorescence labeled sample from a sample tube to the fluorescence detection system. The analyzer can be used for measurements of fluorescent fluids, fluorescence-labeled micro- and nano-particles, and capillary electrophoresis

Key specifications of the sample analyzer include:

- Detecting up to 10,000 objects/s
- Minimized (less than 1%) beads loss and damage
- Sample contacts only with disposable parts of the analyzer (no equipment conservation needed)
- Volumetric speed up to 10 μ l/min
- Sample volume 2-150 μ l
- 1 μ l dead volume
- Various types of beads (spectral range 520nm-720nm)
- Multi-laser excitation

2.2.Operation principle

Schematic of the analyzer is presented in Figure 2.2.1.

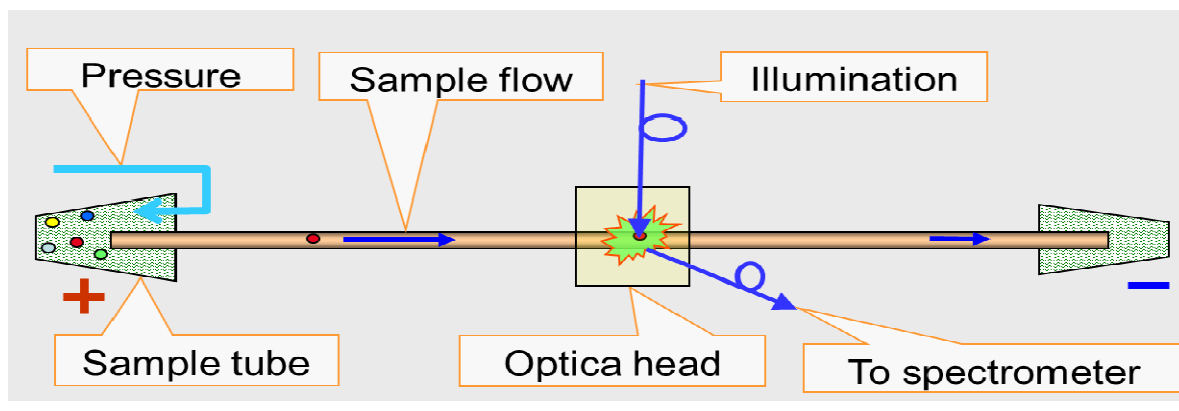


Figure 2.2.1. Sample analyzer

In the sample analyzer, fluorescent objects (molecules, micro-beads) are moved from a sample tube into a capillary by applying either pressure or electric field (electrophoresis). The pressure may be applied by either using fluidic system or compressed air. The capillary is inserted into a fiberized optical head (Figure 2.2.2.).

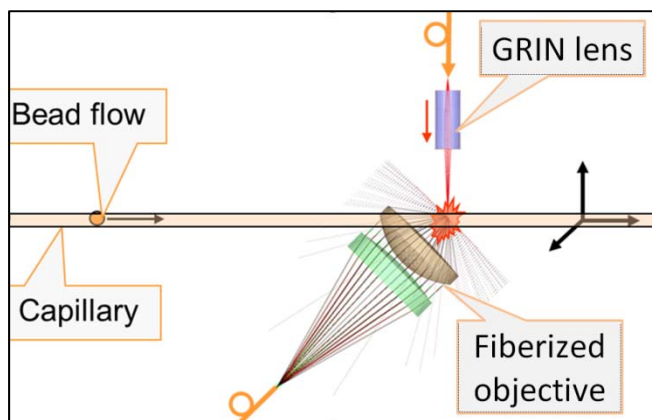


Figure 2.2.2. Schematic of optical head

Illumination from the fiberized laser module is delivered to the optical head and focused through a GRIN lens inside of the capillary. When fluorescent object passes through the laser beam, the excited fluorescence is collected by a fiberized micro-objective and delivered to the fiberized spectrometer, which records fluorescence with single photon sensitivity in 32 spectral

ranges between 500 and 720 nm. The data acquisition speed may be as high as 32MB/s at the frame rate of 1,000,000/s.

2.3. System with fluidic pumping

The purpose of fluidic system is to move samples diluted in liquid substance such as water, various buffers, alcohol, etc. through capillary with constant speed. Design of the fluidic system is shown in Figure 2.3.1. Using 3-way flow switching valve (Upchurch, V-101L) we first move the sample from the sample tube into the pumping syringe (upper panel). From the syringe the sample can be pumped through the capillary (lower panel). When the sample is pumped through the capillary it passes the detection head and fluorescence of the sample is detected. The main advantages of this system are its ability to create high pressure which allows to move objects in the capillary at high speed (up to 1m/s in 25 μ m capillary), and very accurate pressure control, which allows supporting constant pressure (and speed) through the entire measurement process.

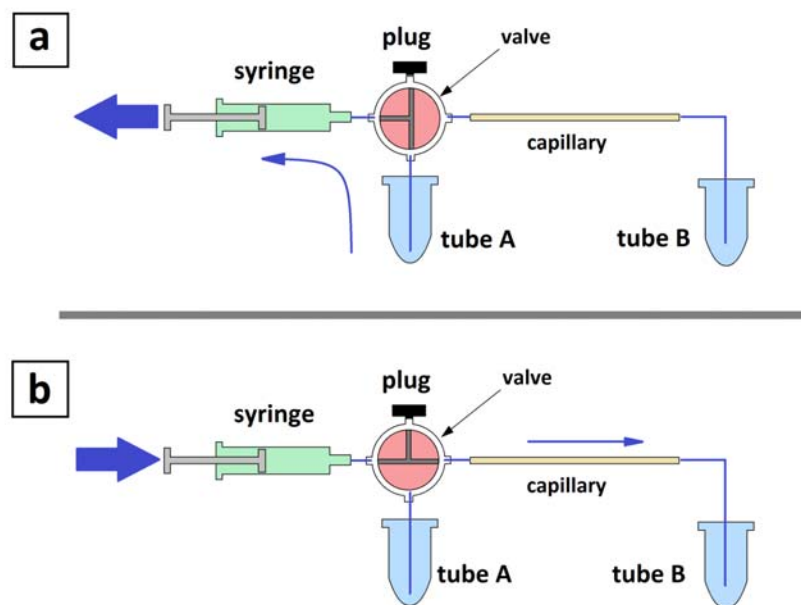


Figure 2.3.1. Schematic of analyzer with fluidic pumping.

In the course of use of the fluidic pumping system we found a number disadvantages when it is used for measurements of micro-objects (beads). The main reason here is that fluids containing beads are moving through a number of valves and tubing connectors. As a result, the

beads can stick to the walls of the tubing, connectors, and inner walls of the valves. This causes a number of major problems. Firstly, the beads which are captured in the valve, damage the valve which ultimately leads to the valve leakage, and also accumulation of the beads in the tubing clogs the system. Second problem is the bead loss itself. Indeed, it is practically impossible to obtain reproducible results because the bead count may change significantly from run to run.

Another disadvantage of the fluidic pumping is that the tubing causes a large dead volume of the system, which requires larger amount of sample.

Therefore our conclusion regarding the fluidic pumping is that this method is good for clear fluids and does not work well with suspensions.

2.4.System with air pressure pumping

In order to avoid disadvantages of the fluidic pumping, I have designed the capillary system with air-pumped sample. The main concept of the system was to avoid any contact of the sample with pumping system. In other words, from the sample tube, the sample should go directly to the capillary.

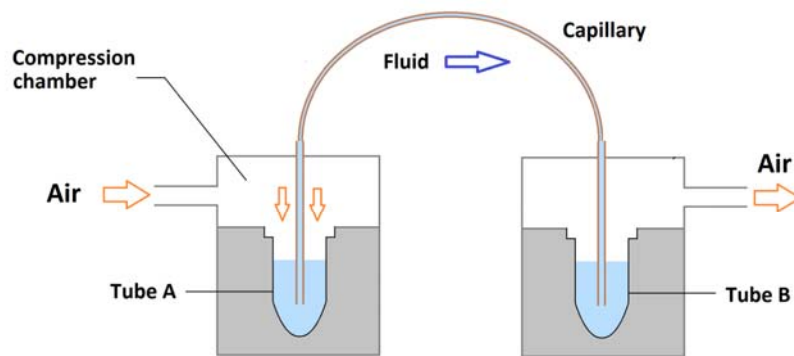


Figure 2.4.1. Schematic of air-pumped analyzer

Schematic of the air-pumped system is shown in Figure 2.4.1.

Sample placed in the tube A and sealed in compression chamber. Compressed air enters into compression chamber and creates pressure on the surface of fluid. Fluid moves through

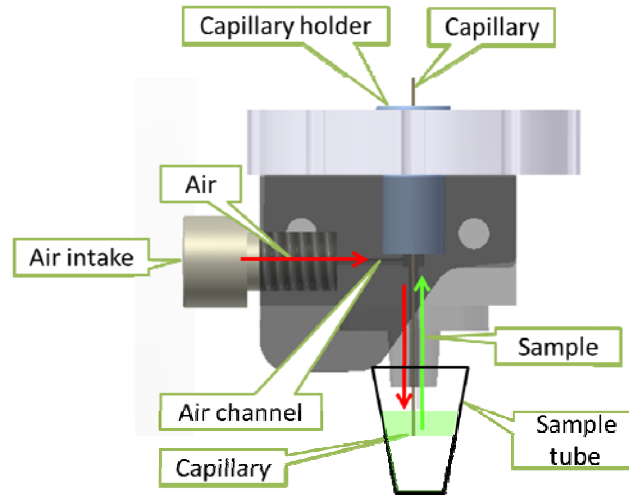


Figure 2.4.2. Detailed design of the compression chamber.

capillary into a waste tube B. Left and right side of the system are symmetric. This design makes possible pushing fluid in both directions by applying air pressure to a left or right input of the fluidic system. It allows, if needed, repeating measurements of the same sample.

The most complex part of the system is the compression chamber. The complexity of this part is caused by (i) a requirement of sealing the sample directly in the standard sample tube, and (ii) also by the necessity of easy replacement of the capillary. A detailed design of the compression chamber is presented in Figure 2.4.2. In the chamber we have T-shape air channel through which air is pumped by a pumping syringe. The capillary is inserted into the chamber and fixed tightly by the capillary holder. Sealing of the chamber volume is done by (i) an airproof connection of the capillary holder with the body of the chamber, and by (ii) pressing the sample tube to the bottom of the chamber. Due to pumping, high air pressure (up to 150psi) is created in the chamber. When the capillary inlet is immersed into the sample tube the sample moves into the capillary due to high air pressure in the chamber.

Figure 2.4.3 presents 3D CAD model and photographs of the manual sample analyzer.

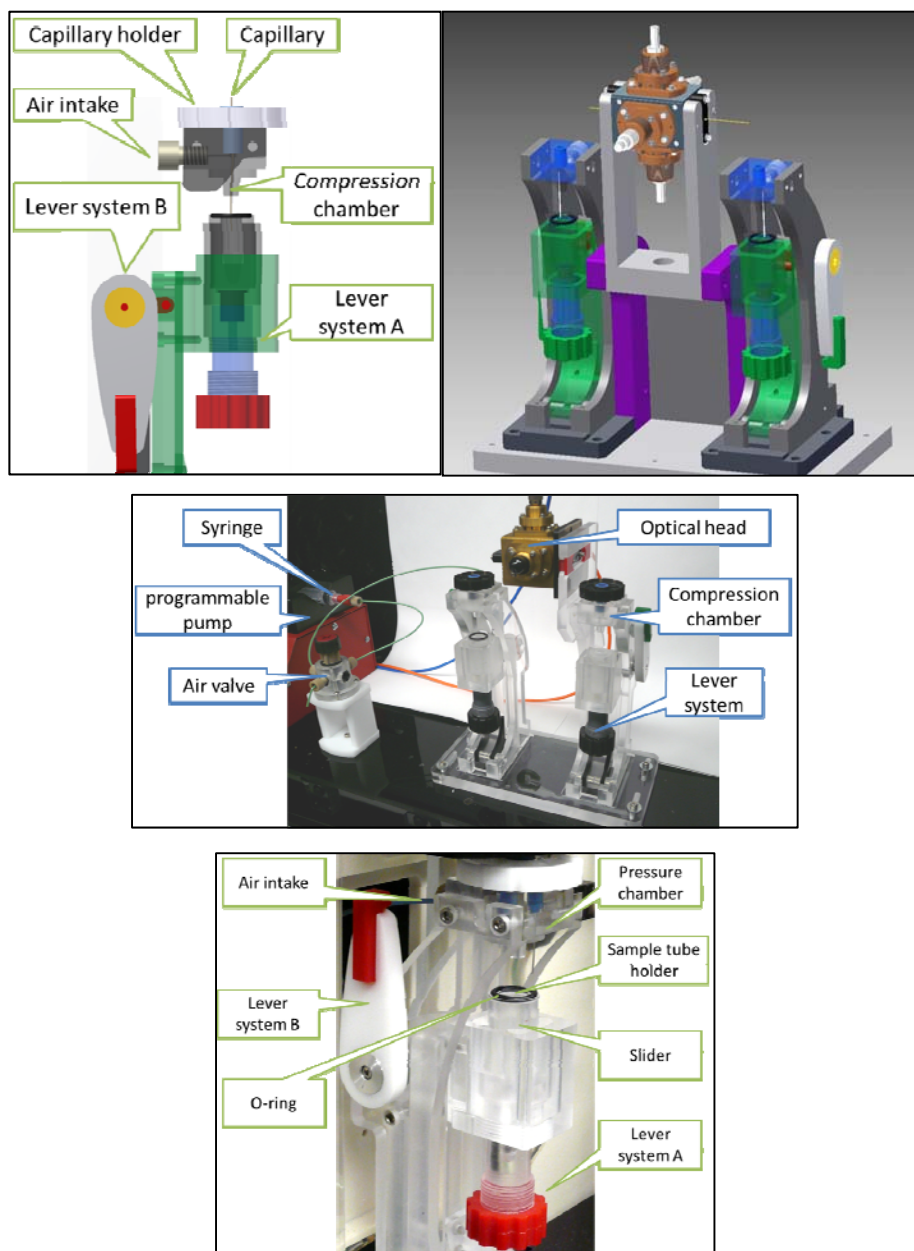


Figure 2.4.3. CAD model (upper panel) and photograph (lower panel) of the air-pumped sample analyzer.

The analyzer comprises two equivalent compression chambers connected by the capillary. The chambers are sealed by pressing the sample tube to the bottom surface of the chamber using the lever system. The capillary is passed through the fiberized optical head (see next section for details). The air is pumped using 5ml syringe mounted on the programmable pump. Using the air valve, pumping can be performed in both directions. The sample tubes are placed on mounts and

using a lever system pressed against the compression chamber, so that the sealed air volume is created.

Because in the system, the analyte goes from the sample tube directly to the capillary, this design has minimum loss of particles during the run. The system design is such that it requires minimum cleaning because all parts that have contact with sample are disposable. Another advantage is minimum dead volume; it means that almost the entire sample volume can be pumped through the capillary.

The use of the syringe pump may cause pressure degradation in cases when sample volume is comparable with the sealed volume of the compression chamber. In this case, the volumetric pumping speed will decrease and this may cause some inaccuracy in measurements. In such cases a compressor which can support the required air pressure must be used.

Some factors add complexity to the design and fabrication of the air-pumped analyzer. Indeed, compression chamber has to be thoroughly sealed, even smallest air leak creates pressure degradation that reduces volumetric speed at which the beads move through the capillary. Another challenge is to design a chamber which is able to hold the required pressure (150psi in our device). Since gases are easily compressible, the volume of the chamber should be as small as possible. Compression chamber also has to be easy re-sealable to enable replacement of sample tubes.

2.5.Fiberized optical head

In the fiberized optical head the collection optics, the illumination optics, and the capillary are positioned along XYZ axes. Figure 2.5.1 presents ray tracing in the optical head.

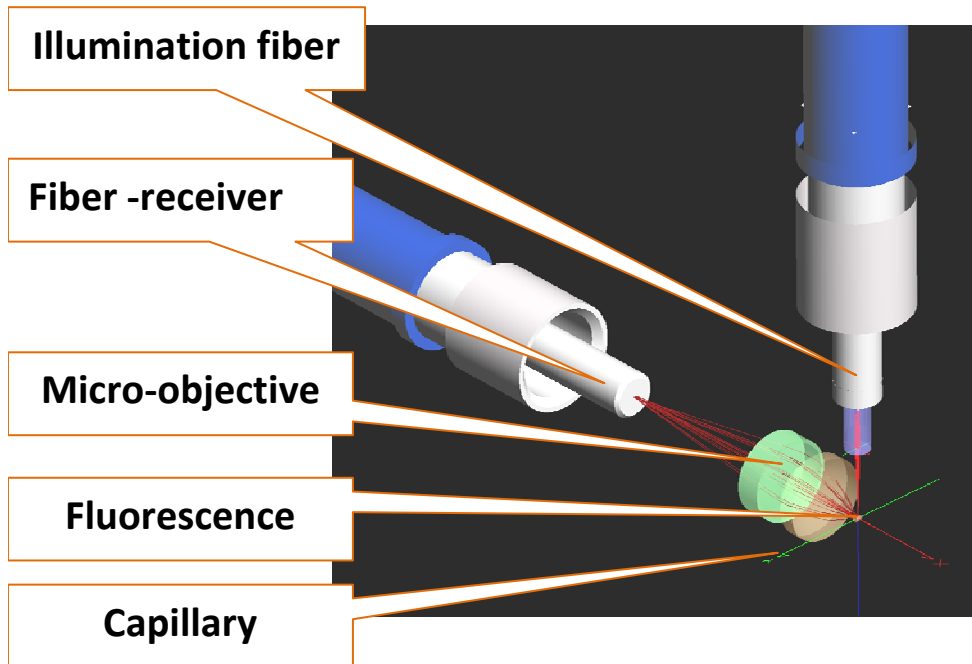


Figure 2.5.1. Ray tracing in optical head

Spectral properties of optical head

Spectral properties of the optical head are very important for accurate decoding of color-encoded beads. Indeed, when the bead is moving through the capillary, it may cross the laser beam along different trajectories (Figure 2.5.2).

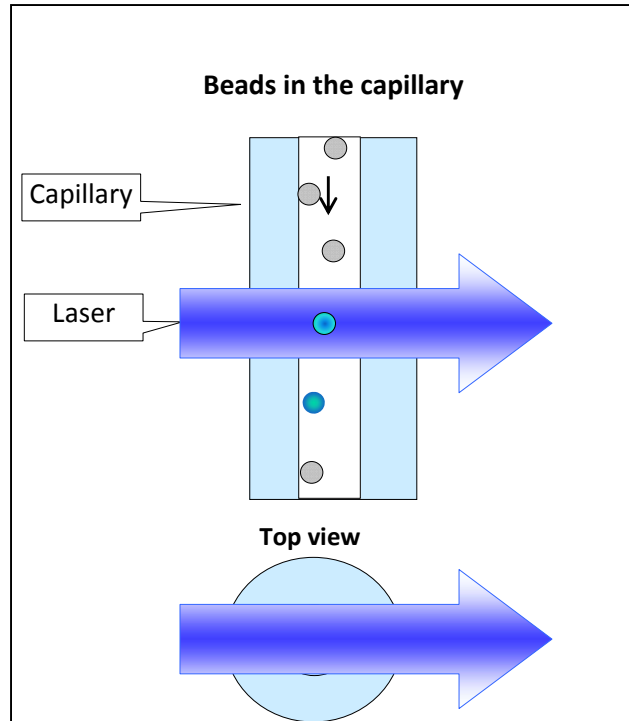


Figure 2.5.2. Beads in the capillary.

Depending on the trajectory, the fluorescence collection efficiency may differ. Moreover, the collection efficiency may depend also on the wavelength of the fluorescence, and these dependences may differ for different wavelengths. Figure 2.5.3 presents schematic of mutual positions of the capillary and collection optics in optical head (the capillary goes along Z axis).

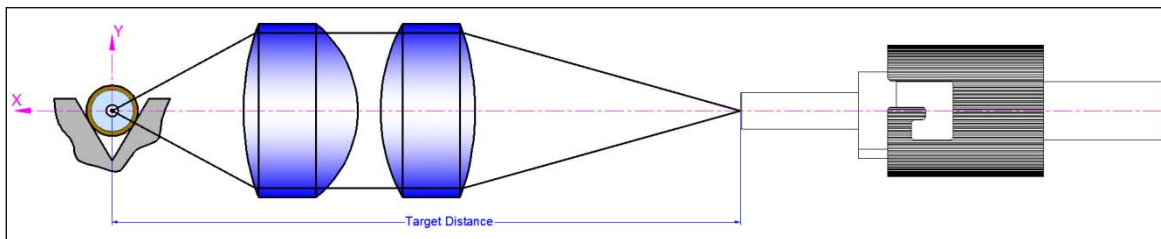


Figure 2.5.3. Mutual positions of the capillary, collection optics, and fiber-receiver in the optical head.

Upper and middle panels of Figure 2.5.4 show results of ray tracing for three wavelengths (520nm, 620nm, and 720nm) for the bead positioned in the center of the capillary cross section ($X=0$, $Y=0$).

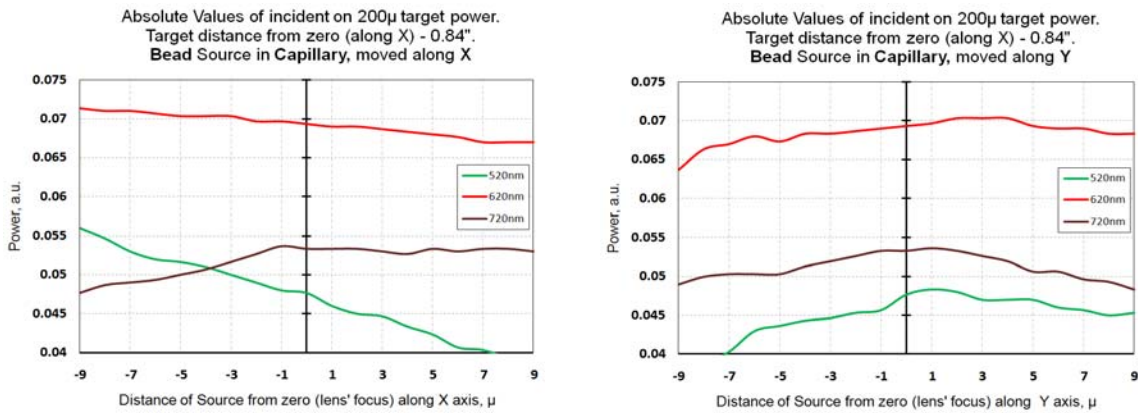
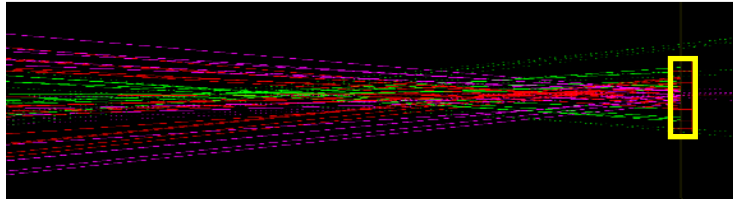
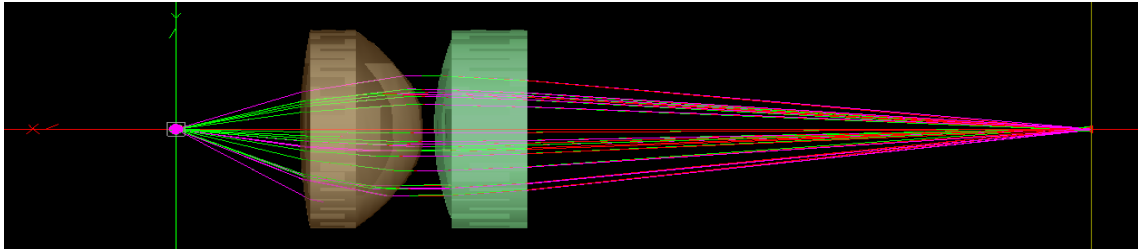


Figure.2.5.4. Simulation of fluorescence collection efficiency for bead in capillary. (200 μ collection fiber). Full ray tracing (upper panel), enlarged ray intercept area (middle panel), collection efficiency (lower panel).

As can be seen, focus distances differ for all three wavelengths (middle panel). This is why collection efficiency of the fiber-receiver will depend on the wavelength. When fluorescent beads go along the capillary, their XY coordinates vary within diameter of the capillary. Obviously, position of the bead in the capillary influences its focusing in the collection system. This is why fluorescence collection efficiency depends on both wavelength and position of the bead (Figure 2.5.4, lower panel).

The presented situation tells that the decoding of color-encoded beads will depend on trajectories of the beads in the capillary. As can be determined from calculations, the inaccuracy in determination of color code may be as high as 10-15%.

Figure 2.5.5 presents results of simulation of fluorescence collection efficiency in 200 μ collection fiber using achromatic collection lenses. As can be seen, from the ray tracing, with achromatic objective all wavelengths are focused in the same position (upper and middle panels).

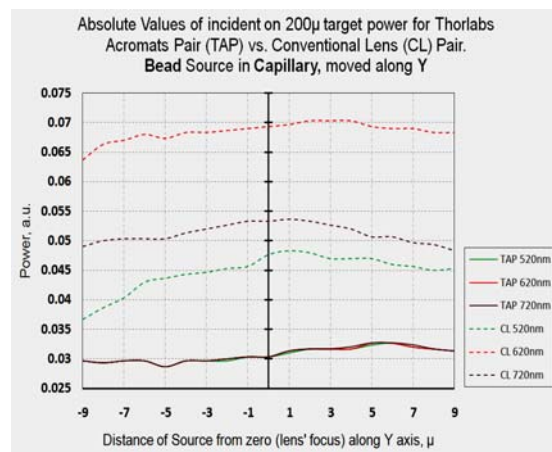
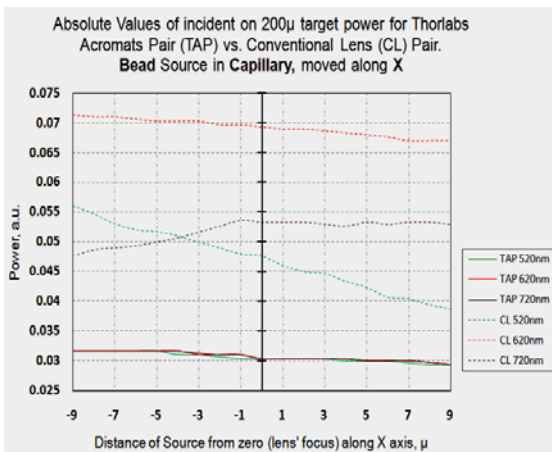
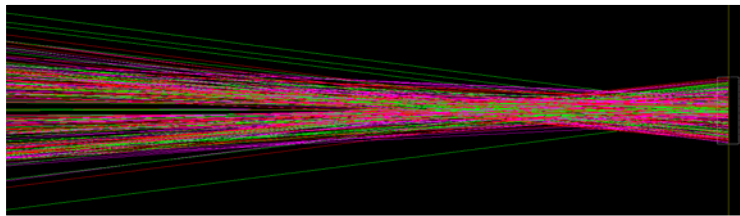
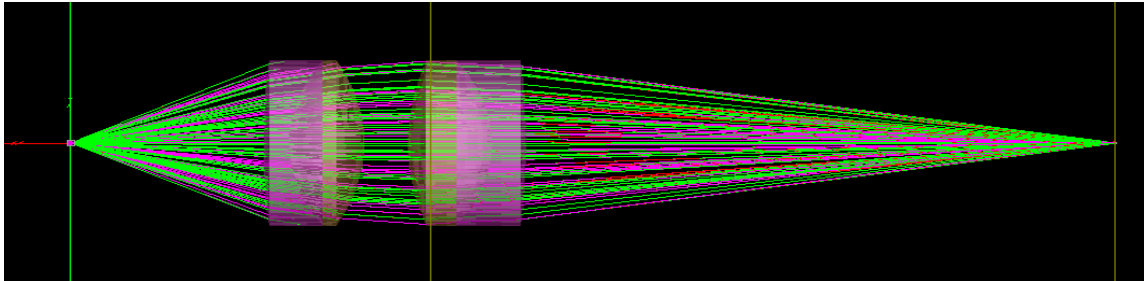
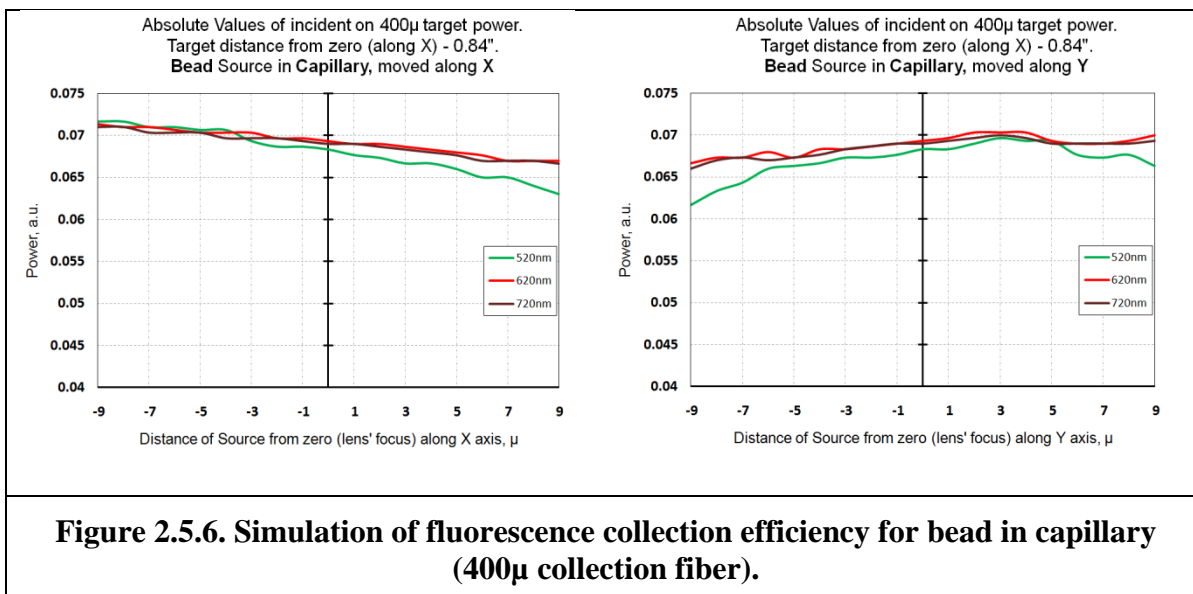


Figure 2.5.5. Simulation of fluorescence collection efficiency for bead in capillary. (achromatic collection lenses, 200 μ collection fiber). Full ray tracing (upper panel), enlarged ray intercept area (middle panel), collection efficiency (lower panel).

Fluorescence collection efficiencies do not depend on wavelengths, but the use of achromatic optics yields absolute collection efficiency 2.5-3 fold lower than in the previous case (lower panel).

Our calculations also show that dependence of collection efficiency on fluorescence wavelength varies for different diameter of collection fiber (Figure 2.5.6).



However, the use of 400µm collection fiber is undesirable because it reduces spectral resolution of the spectrometer very significantly.

2.6. Testing of the bead reader

Volumetric speed depending on:

- Capillary diameter – 25µm (20cm lengths) speed in water is 2ml/min
- Capillary diameter – 50µm (20cm lengths) speed in water is 30ml/min
- Capillary diameter – 50µm (55cm lengths) speed in POP-7 gel is 8mm/min

Loss of beads less than 1%

Characterization detection accuracy

We have carried out characterization of the accuracy and the reproducibility of measured ratios of fluorescent signals obtained from groups of identical beads labeled with two dyes.

In order to check the accuracy of bead detection, we have assembled a special measurement setup (Figure 2.6.1 upper panel), which allowed us to control lateral position of the capillary in the optical head as well as to rotate the capillary around its axis. Using a high precision positioning system, the capillary was moved so that beads were illuminated individually by the laser beam, and multiple measurements of the same bead in a fixed position were performed and

recorded. Figure 2.6.1 (middle panel) shows a distribution of photon counts measured from an individual bead (see also a red area in Figure. 2.6.1 (lower panel)). As can be seen, the

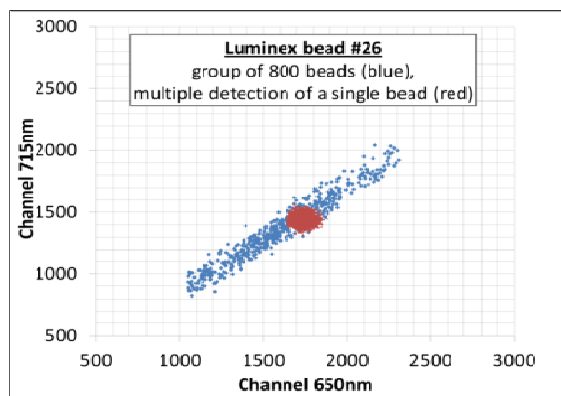
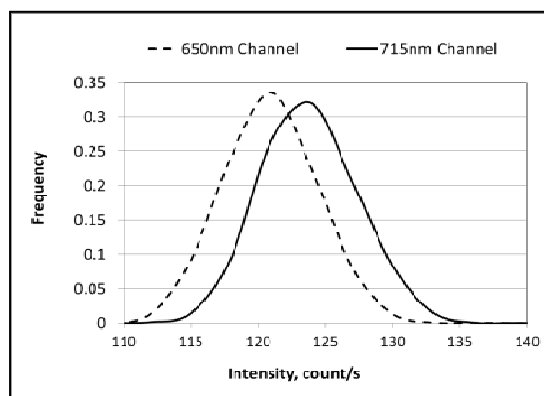
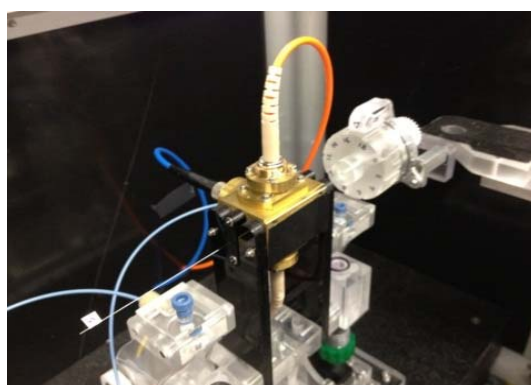


Figure 2.6.1. Measurement setup for determination of bead detection accuracy (upper panel), distribution of photon counts measured from individual bead immobilized in the optical head (middle panel), and a bead map of a set of Luminex beads (lower panel).

distribution is Poissonian (*mean is equal to variance*), which is characteristic of photon counting. Thus, we conclude that the accuracy of our measurement system is limited by photon counting noise, which is a fundamental limit. Figure 2.6.1 (lower panel) presents a bead map obtained for a group of ~800 Luminex beads (blue dots) in comparison with multiple measurements of a single bead, which was immobilized in the optical head (red dots).

A large (~2-fold) spread of measured photon counts from individual beads (blue dots) is due to a non-uniform distribution of the laser beam in the capillary. However, the ratio of photon counts obtained in two spectral channels is almost constant, and it varies mostly due to the photon counting noise (see red zone indicating single bead measurements).

3. CHAPTER 3. APPLICATIONS

3.1.Measurement of composition of fluorescent dyes

Figure.3.1.1 presents a schematic of the setup for measurements of fluorescence spectra of multicolor fluorescent analytes in capillary. In the setup, the analyte is moved from a sample tube into the capillary by either electrophoresis or air pressure. The capillary is inserted into a fiberized optical system. Fluorescence in the analyte is excited by Ar-ion laser (488 and 514 nm, Uniphase, CA, USA), collected by the fiber (200 μ core diameter) and delivered to the spectral separation module of the spectrometer.

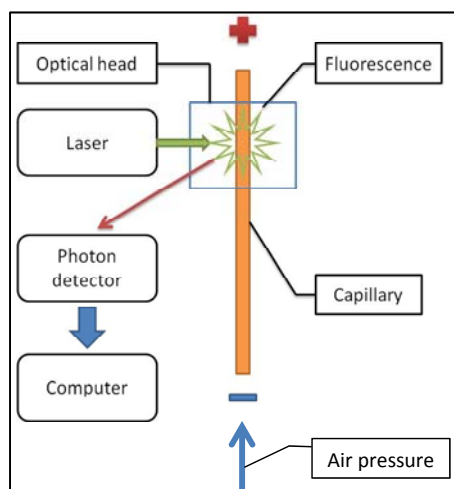


Figure.3.1.1. Schematic of the measurement setup.

Our task in this application is to determine composition of several fluorescent dyes which are mixed in a single sample. In order to do this we measure fluorescent spectra of individual dyes using our measurement setup, and then use known spectra for determination of content of individual dyes in the analyzed mixture.

The experimental setup records optical fluorescence spectra in 32 independent channels of the PMT. Each sample represents the full spectrum of fluorescence in the range from 460 to 630 nm. The range is divided among the PMT channels so that each channel collects fluorescence in a 5-10 nm band.

Fluorescence from an individual dye is simultaneously measured by several PMT channels. Photocount rate is distributed among PMT channels according to fluorescence spectrum of the dye. In practice, the measured distribution of photocount rates is influenced by differences in sensitivity between channels and optical cross-talk. In order to ensure reliable operation, the system must be calibrated for a particular dye set. During the calibration some unit concentration of each dye of the set is presented to the system and its spectrum is recorded. The set of spectra is used to estimate optimal parameters of signal processing (see [Tsupryk A., 2008]). Detailed description of color decomposition method is also presented in [D. Gudkov et al, 2012, submitted to Biosensors & Bioelectronics].

Measurements of mixtures of quantum dots

In order to test the sample analyzer, we have carried out measurements of mixtures of three different types of QDs with strong spectral overlap (Figure 3.1.2). QDs re-suspended in toluene 1mg/ml were obtained from Crystalplex, Pittsburgh PA.

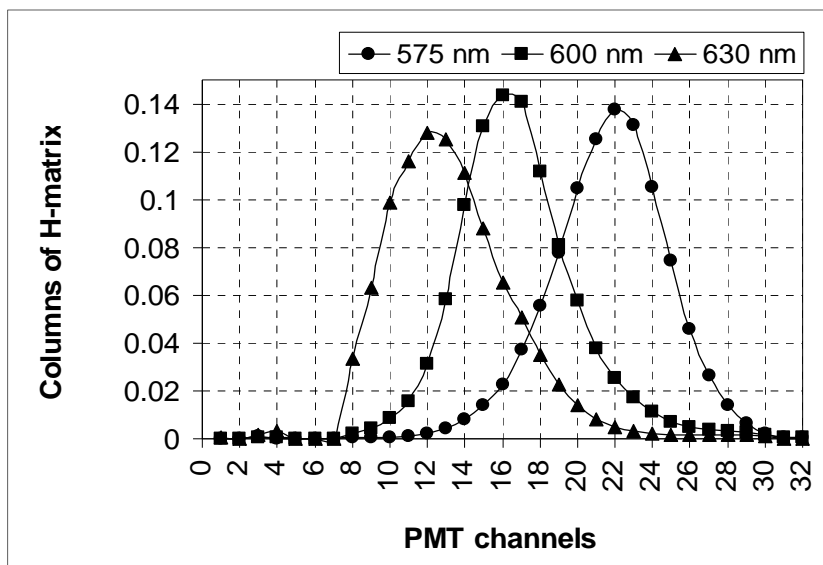


Figure.3.1.2. Spectra of individual quantum dots.

We premixed QD solution consisting of 50% of 575 nm QDs and 50% of 630 nm QDs and prepared a subsequent series of mixtures in which the content of 600 nm QDs varied from 0% to 15%. Mixtures of quantum dots were pumped into fused silica capillaries with transparent coating (360 μm OD, 100 μm ID, Polymicro AZ, USA) and detected using the setup presented in

Figure 3.1.1. The total recording time for each mixture was about 5 minutes with minimum integration time of ~ 20 ms which provided approximately 15,000 data points per mixture. The color deconvolution procedure described D. Gudkov, 2012, was applied to photo-counts obtained in 32 channels at each data point during 20 ms photon integration time. In order to increase the photon integration time, we summed photo-counts obtained from several sequential data points and thus applied the color deconvolution procedure to longer time intervals.

Based on the measured spectra of individual QDs we have carried out color deconvolution procedure. The obtained results are presented in Figure 3.1.3. As can be seen from Figure 3.1.3 (left), the concentration of the 600 nm QDs determined from the fluorescence measurements is practically equal to the original concentration of that component in the prepared QD mixtures. Fig.3.1.3 (right) presents distribution of concentration of the 600 nm QDs recovered from multiple measurements of the mixtures with 0.1 s and 0.02 s integration times. As can be seen, increase of the integration time improves the resolution of QD mixtures.

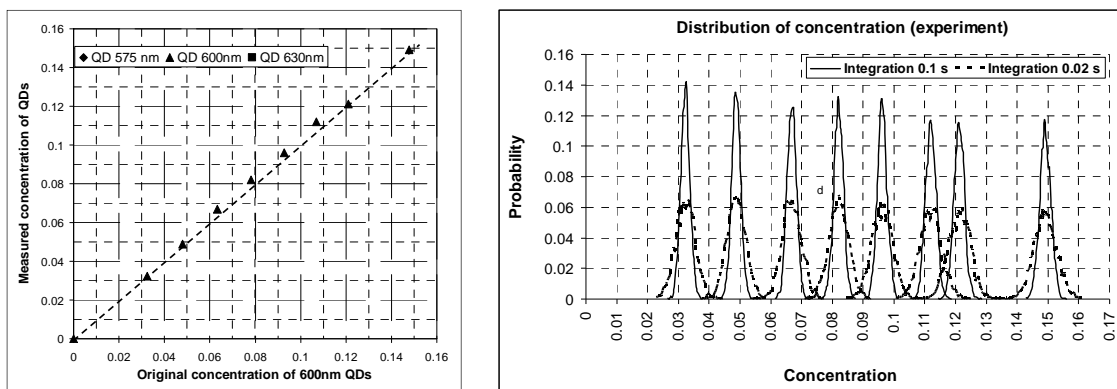


Figure 3.1.3. Deconvolution of QD mixtures (left) and distributions of concentrations of 600 nm QD determined with 0.1 s and 0.02 s integration times (right).

Based on the analysis of distributions of color compositions obtained experimentally found that distribution of measurements is Poissonian, and the accuracy is determined by square root of the obtained photon count.

As can be found from Figure 3.1.3, the developed sample analyzer can resolve dye mixture with precision of $\sim 3\%$, which is very important for measurements of color-encoded objects (see next section).

3.2.Measurement of color encoded beads

Beads encoded with combinations of up to 5 types of quantum dots

Characterization of QDs

Five types of QDs (525nm, 540nm, 580nm, 630nm and 665nm) were purchased from CrystalPlex Corp, Pittsburgh, PA. Schematic of the setup used for characterization of QDs is presented in Figure.3.1.1. Illumination of the capillary was done by single mode lasers [*Ar-ion*, (488nm and 514nm), *solid state* (532nm), and *semiconductor* (638nm)]. We found that the excitation by Ar-ion laser (25 μ m beam, 25mW output power) was optimal for all QD types. Stock solutions of QDs (8 μ M) were diluted 1:5000 and run through 50 μ m ID capillary. The obtained QD spectra are shown in Fig.3.2.1.

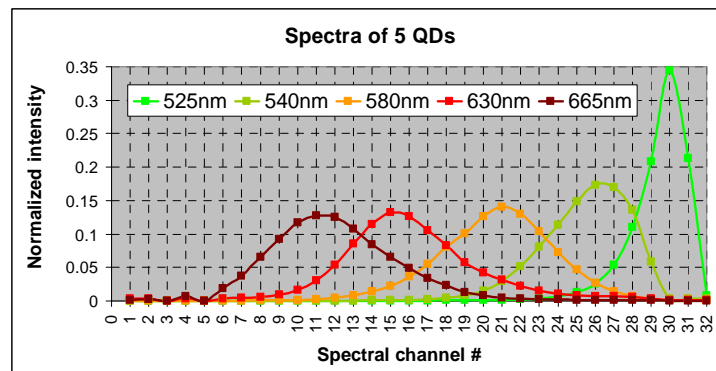


Figure. 3.2.1. Emission spectra of five QDs.

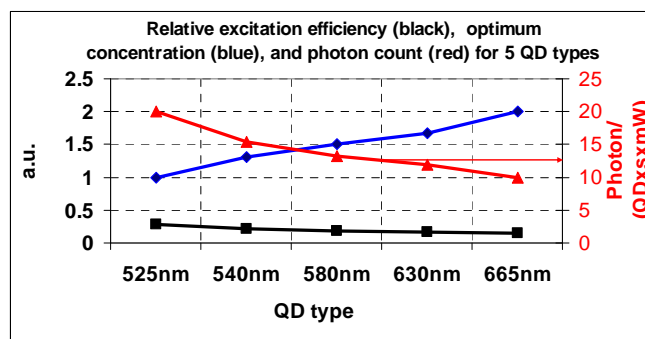


Figure 3.2.2. Excitation efficiency, optimum concentration and photon count for 5 QDs of five QDs.

For each QD type we have measured the photon yield (*a number of photons emitted by one QD in the maximum spectral channel per second per 1mW of illumination power*, Figure.3.2.2, red

curve). Based on the relative excitation efficiency of all studied QD types (Figure 3.2.2, black curve), we have determined the optimum relative concentrations of QDs which provide equal photocount rate for all five QD types (Figure 3.2.2, blue curve). All calculations were done based on the illuminated volume of the capillary (~31,000 cubic microns), concentration of QDs in the solution, and the illumination power (25mW from Ar-ion laser).

Characterization of beads 5µm, single color porous beads doped with 525nm, 540nm, 580nm, 630nm, and 665nm QDs were purchased from CrystalPlex Corp, Pittsburgh, PA.

Imaging of the beads showed that the variation of beads' diameter did not exceed 20%.

Saturation of the beads According to the manufacturer, the number of QDs which saturates the bead can be as high as ~2million. Based on the determined excitation efficiencies for of QDs and targeting equal fluorescence signal from each QD type, we have calculated a maximum number of QDs per bead for sets of beads encoded with 2-, 3-, 4- and 5 QD types (Figure 3.2.3.).

It is important to emphasize, that the number of QDs per bead may be limited by the linearity of the photon detector (10^8 photon/s in our spectrometer) rather than by the saturation of the bead by QDs. Our estimates based on the measured photon yield of QDs show that for 25mW excitation power the number of QDs per bead may not exceed ~250,000.

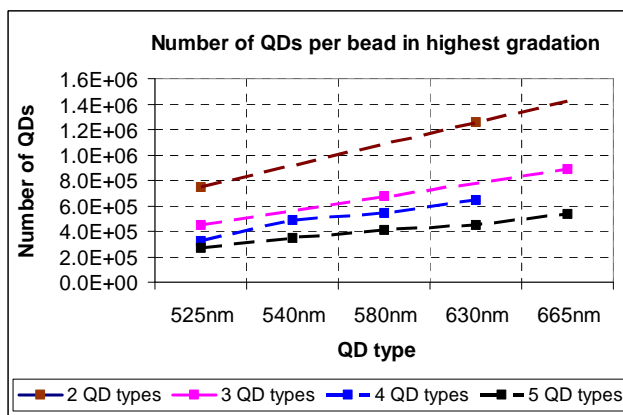


Figure 3.2.3. Emission spectra of five QDs.

Testing statistically encoded beads

According to the determined specifications, 2-, 3-, 4-, and 5- color sets of randomly encoded beads were fabricated for us by CrystalPlex Corp.

In order to assess quality of the random coloring technology, we introduced two major characteristics for the set of beads: **(i)** a number of distinguishable codes in the set, and **(ii)** distribution of random codes on the color map (*in the phase space of color codes*).

(i) Determination of distinguishable codes

We call a bead unique (having a distinguishable code), if its color code differs *significantly* from codes of all other beads in the set. In this case, once detected, the unique bead becomes identifiable by its color code in all subsequent detection events.

Criterion for recognition of the beads and assigning accuracy to the bead decoding In order to distinguish two statistically encoded beads, we introduce a parameter called “code distance” $D_{12} = \sqrt{\sum_{i=1}^{N_{QDS}} (C_{i1} - C_{i2})^2}$, where C_{i1} and C_{i2} are color codes of the two beads for i^{th} color component and N_{QDS} is a number of QD types used for encoding the beads.

In single photon spectrometer, the accuracy of the determination of the color code of the bead is limited by Poisson noise generated by fluorescence photons. Therefore, 99% of all measurements fall within the interval $\delta_{TH} = \pm 3\sigma$ where σ is standard deviation, which is related to the detected number of photons Φ as: $\sigma = \sqrt{\Phi}$. We consider codes of two beads distinguishable if the code distance $D_{1,2} > \delta_{TH}$ (1).

In order to determine experimentally parameter δ_{TH} we carried out multiple measurements of the same bead by fixing position of the bead in the capillary using liquid polymer matrix. We performed spectral measurements of the bead (Figure 3.2.4, upper panel) and calculated the distances between codes obtained from all measurements. In order to determine δ_{TH} we have developed specialized software which calculates code distances between all beads in the set. Figure.3.2.4, middle panel, presents a distribution of normalized code distances calculated for 4,000 measurements of a single bead. As can be seen, distances between the measured codes of the bead do not exceed 0.003 (which means that if the code distance between two beads exceeds 0.003, the beads are distinguishable with 99% accuracy).

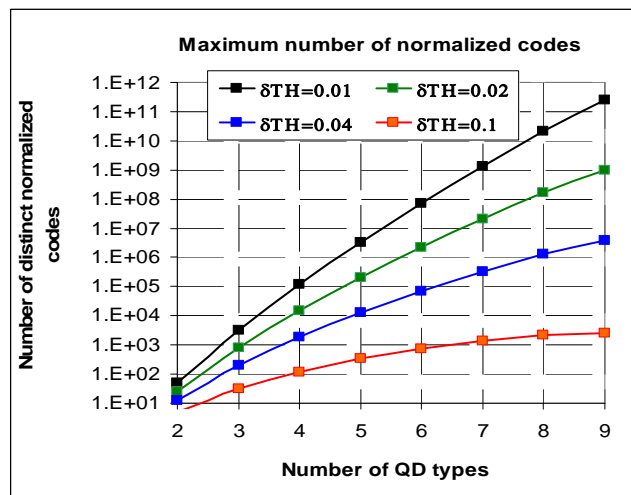
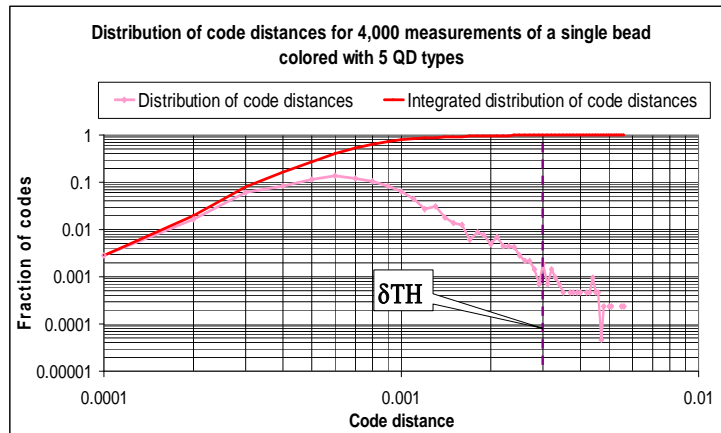
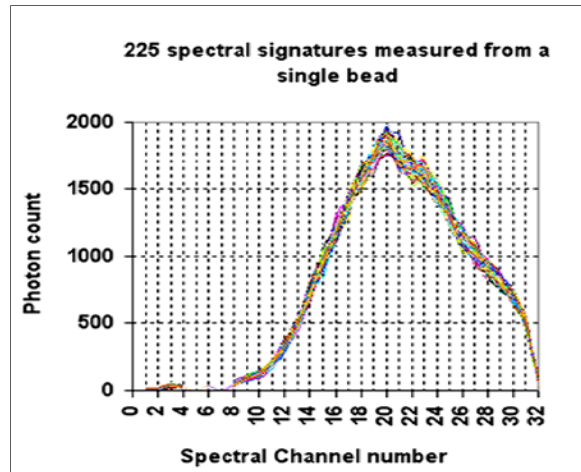


Figure 3.2.4. Spectra of a single bead (upper panel), determination of δ_{TH} (middle panel), number of distinct color combination (lower panel).

Statistics of random codes, distribution of codes on color map

Number of distinguishable color codes

Using Monte-Carlo technique, we have calculated a number of distinguishable normalized color codes in sets of beads depending on the number of encoding QD types and δ_{TH} as a parameter. Results of computer simulations of the normalized color codes for randomly encoded sets of beads are presented in Figure 3.2.4, lower panel (components of a normalized color code present fractions of photon counts generated by each type of encoding QDs. Sum of all code components is always equal to 1).

Based on the obtained results, we conclude that for $\delta_{TH} > 0.003$ one can potentially obtain up to 100,000 distinct color codes by using only 4 QD types and more than a million codes by using 5 QD types.

2-color sets of beads

Using the developed sample analyzer measurements of four 2-color random sets of beads: (525/540nm, 600/630nm, 540/580nm and 660/660nm) have been carried out and yielded similar results for all four sets. Based on the obtained δ_{TH} , we have found up to 40-50 distinguishable codes in the measured 2-color sets of beads. Typical color maps for absolute and normalized codes for 2-color set are presented in Figure 3.2.5 (left and middle).

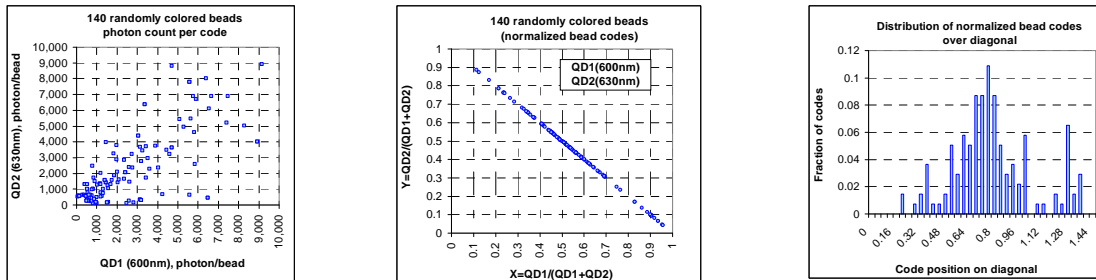


Figure 3.2.5. Statistics of color codes in 2-color random set (600nm/630nm).

Since the sum of normalized codes for any individual bead is always equal to 1, the color map for the normalized codes is a diagonal of the square (1, 1). The distribution of the normalized codes over the diagonal (Figure 3.2.5, right) shows that the color map is filled with

the codes rather unevenly. It means that in order to provide maximum number of codes predicted in Figure 3.2.5, *right panel*, the coloring technology needs to be optimized, targeting more even distribution of the codes.

3-, 4-, and 5-color sets of beads

Typical distributions of codes measured for random sets of beads encoded with two and more QD types are shown on radial color maps (Figure 3.2.6, *upper panels*). Each axis of the radial map presents one component of the normalized code. Individual beads are shown on the map as polygons whose sides connect values of the beads' codes on the axes of the map.

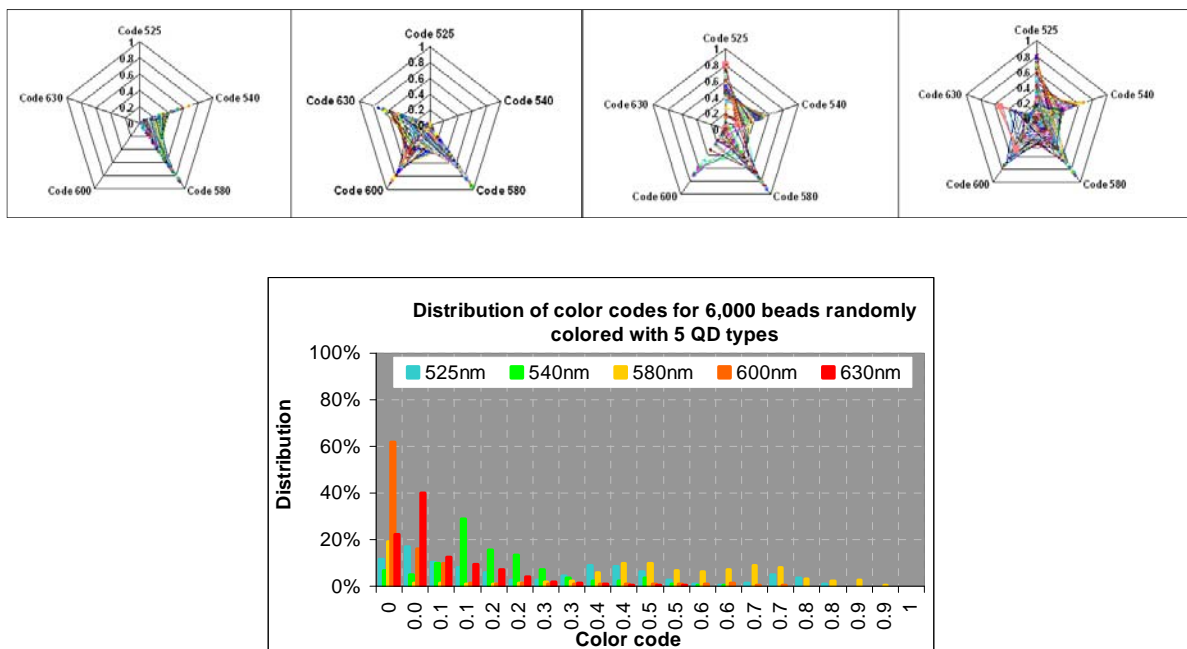


Figure 3.2.6. Upper panels: typical color maps for 2-, 3-, 4-, and 5-color random sets of beads. Lower panel: distribution of color codes experimentally obtained for 5-color set of 6,000 beads.

Lower panel of Figure 3.2.6 shows a distribution of color codes measured for 5-color sets. This shape is typical for all studies sets of beads. As can be seen from the obtained data, the space on the color maps occupied by the codes of the beads covers only a small fraction of the map's area. Because of that, the number of distinguishable codes, obtained for the sets is smaller than maximum number of codes, which can be reached only if the entire color map is filled evenly with random codes.

Taking into account the experimentally determined δ_{TH} , we have been able to identify up to 40-50 distinguishable codes with 2-color sets; up to 150 codes with 3-color sets; up to 1,000 codes with 4-color sets; and more than 9,000 codes with 5-color sets.

Detection of color encoded beads from Luminex

In order to evaluate the accuracy of our system in the bead recognition, we have carried out experiments with commercial beads from Luminex (Microplex X-Tag Microspheres, Luminex, Austin TX, USA) which are internally labeled with 630 nm and 715 nm dyes having certain intensity levels of the two dyes. Figure 3.2.7 shows spectra of two encoding dyes (left) and the bead-map obtained for 50 types of beads in 50 μm ID capillary (right).

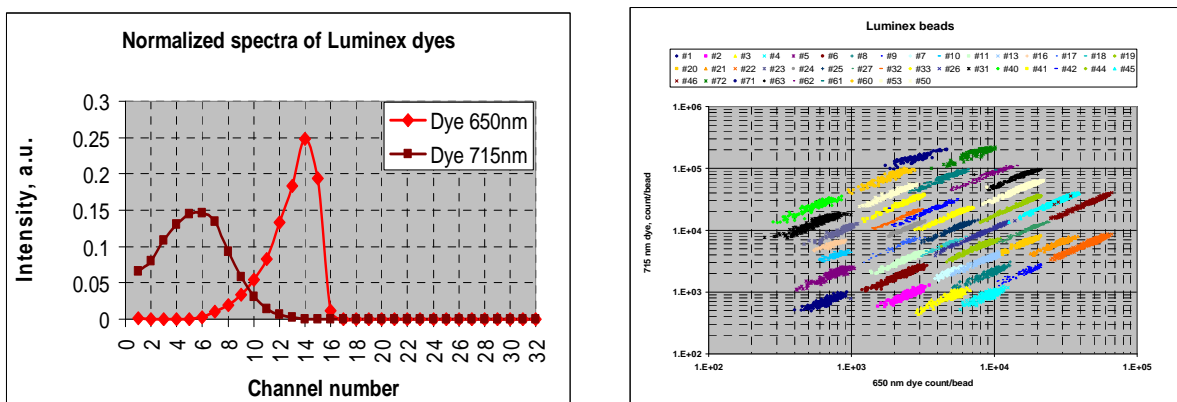


Figure 3.2.7. Spectra of Luminex dyes and bead-map of 50 types of Luminex beads.

As can be seen, the areas occupied by different types of beads do not overlap. Accuracy of the code recognition is determined as a ratio of the correctly classified beads to the total number of the detected beads. Our data analysis showed that we could recognize up to 50 types of beads with 99% accuracy.

Typical color maps for two different sets of beads are shown in the Figure 3.2.8 (*X and Y axes present photon count obtained from individual beads from 650 nm and 715 nm dyes respectively*).

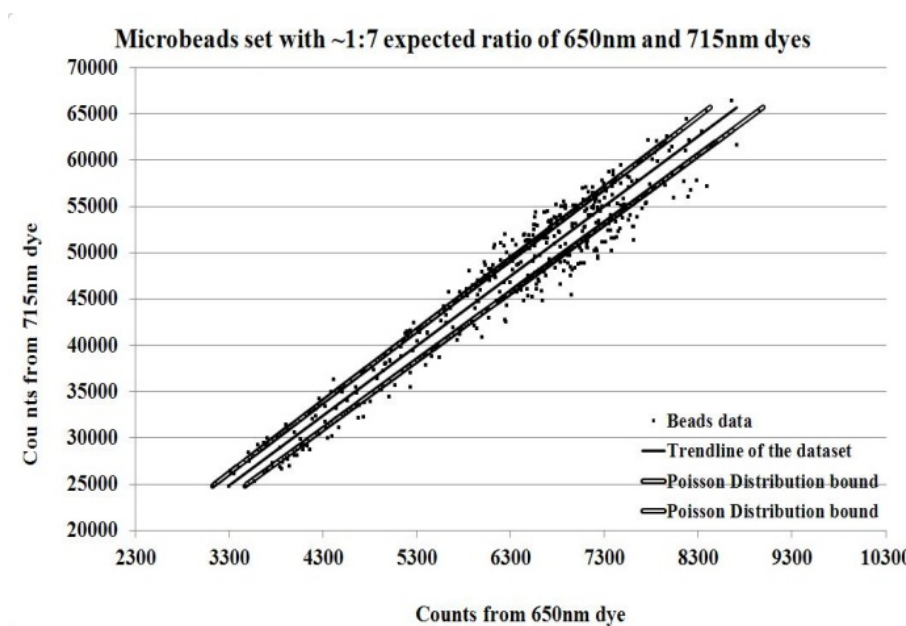
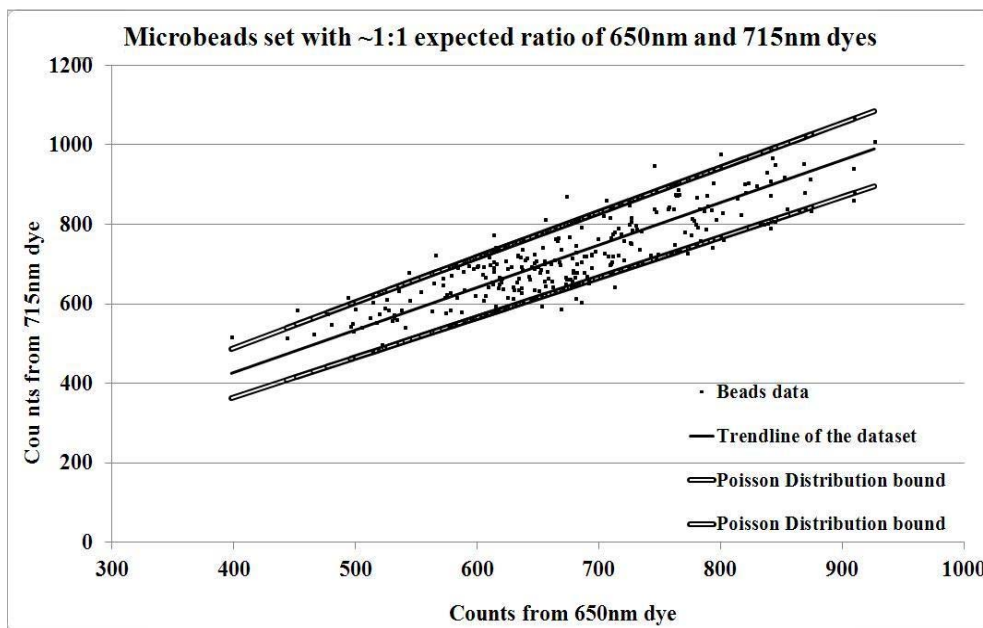


Figure 3.2.8. Color maps obtained from two types of Luminex beads with two different ratios of 630 nm and 715 nm dyes. Upper and lower panels present 1:1 and 1:7 ratios, respectively.

As can be seen, the beads of the same type are distributed along both axes which can be attributed to a non-uniform distribution of the laser power inside the capillary. Trend line represents the slope corresponding to the ratio of the two dyes. Double lines represent 3 standard ranges of deviation (3σ) from corresponding photon counts on the trend line. For both sets of

beads one can see significant number of beads detected out of 3σ range, which in our opinion shows a non-uniformity in the bead coloring process.

3.3.DNA sequencing

DNA sequencing system

The developed sample analyzer was tested for DNA sequencing by capillary electrophoresis. Photograph and schematic of DNA sequencing system are presented in Fig.3.3.1. In order to perform DNA sequencing, 50cm, 50 μ m ID capillary was filled with liquid polymer. In order to move DNA sample through the capillary, high voltage (8-15kV) was

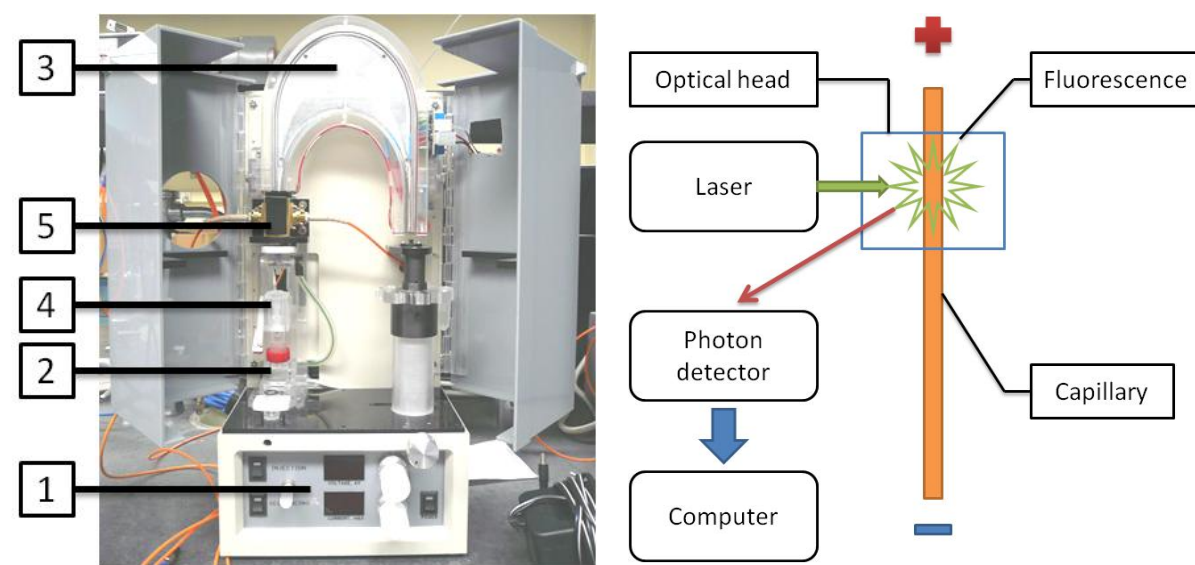


Figure 3.3.1. Photograph and schematic of the single lane DNA sequencer.

Right panel: sample analyzer modified for DNA sequencing. 1 - high voltage supply with a built-in voltmeter and ampermeter; 2 - polymer replacement system; 3 - temperature control system; 4 - tube-changer; 5 – fiberized precision optical system.

Left panel: schematic of the optical system.

applied to the sample and the waste tubes. When the fluorescently labeled DNA fragments pass the optical system (B), they are illuminated by a fiberized laser (either Ar-ion laser (488 and 514nm, 20mW, Uniphase, CA, USA) or Nd:YAG laser (532nm, 25mW, GCL-025-S,

CrystaLaser, NV, USA). The laser-induced fluorescence is collected by the fiber (200µm core diameter) and delivered to the spectral separation module.

DNA sequencing chemistry and protocols

DNA samples labeled with BigDye (BigDye Terminator v1.1 Sequencing Standard) obtained from Applied Biosystems (Foster City, CA, USA) were denatured in 25 µl Formamide (Applied Biosystems) followed by dilution 1 to 5 in HPLC grade water immediately prior to injection. DNA separation was conducted at 150 V/cm in 40 cm (35 cm separation length) uncoated capillaries (50 µm inner diameter, 365 µm outer diameter from Polymicro, Phoenix, AZ, USA) filled with POP-7 separation medium (Applied Biosystems) at 50°C. Applied Biosystems running buffer was used.

Processing of DNA sequencing data

We tested two approaches to estimation of fluorescence intensities produced by four dyes used for labeling of A, C, G and T terminators: decomposition of the entire measured spectrum

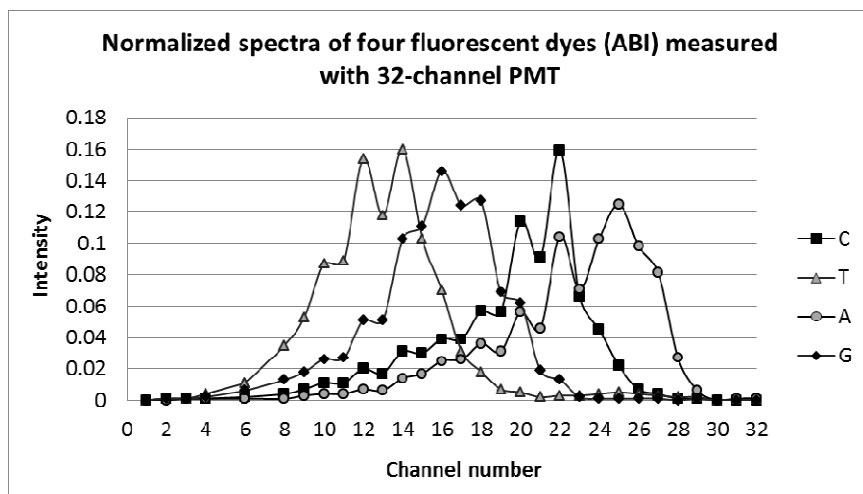


Figure 3.3.2. Spectra of labeled terminators measured by 32-channel spectrometer.

and application of virtual filters (see Tsupryk A. 2008 for signal processing details). For the *decomposition of the entire measured spectrum*, the spectra of individual dyes were obtained

experimentally by measuring fluorescence emitted by each dye separately and by normalizing the measured vector of photocount rates (Figure 3.3.2).

The determination of concentrations of individual dyes in the mixture was performed for each time frame. The obtained sequence of dye concentrations forms four sequencing traces, which are further used for base calling.

The *method of virtual filters* is a simplified approach, which can be successfully used if the dye spectra are well separated. Four virtual filters are formed by selecting a certain PMT channels or groups of channels so that fluorescence of each dye contributes the most into measurements in the channel, which was assigned to it, and minimum response in the other channels. For the simplified data processing we selected 18, 15, 11 and 5 channels of the PMT for G, A, T and C correspondingly.

The obtained color de-convolution matrices for Ar-ion and Nd:YAG lasers are shown in Figure 3.3.3.

Excitation	Ar-ion laser (488nm & 514nm), 18mW			
Base	Color deconvolution matrix			
G	1	0.842	0.297	0.124
A	0.147	1	0.366	0.264
T	0.007	0.104	1	0.809
C	0	0.008	0.057	1
Excitation	Nd-YAG laser (532nm), 17mW			
Base	Color deconvolution matrix			
G	1	0.679	0.254	0.084
A	0.416	1	0.285	0.179
T	0.009	0.253	1	0.493
C	0.025	0.024	0.109	1

Figure 3.3.3. System matrix C and color deconvolution matrices used for processing of DNA sequencing data.

In both data processing approaches the obtained four sequencing traces undergo standard processing [Alaverdyan L., 2002].

It is important to emphasize that both approaches to the data processing yielded similar sequencing traces and base calling quality. However, for measurements characterized by small signal-to-noise ratios the first approach gives better results since it uses the entire information about the fluorescence emitted by each dye rather than just the fluorescence obtained in only four selected channels of the spectrometer (see Figure 3.3.4).

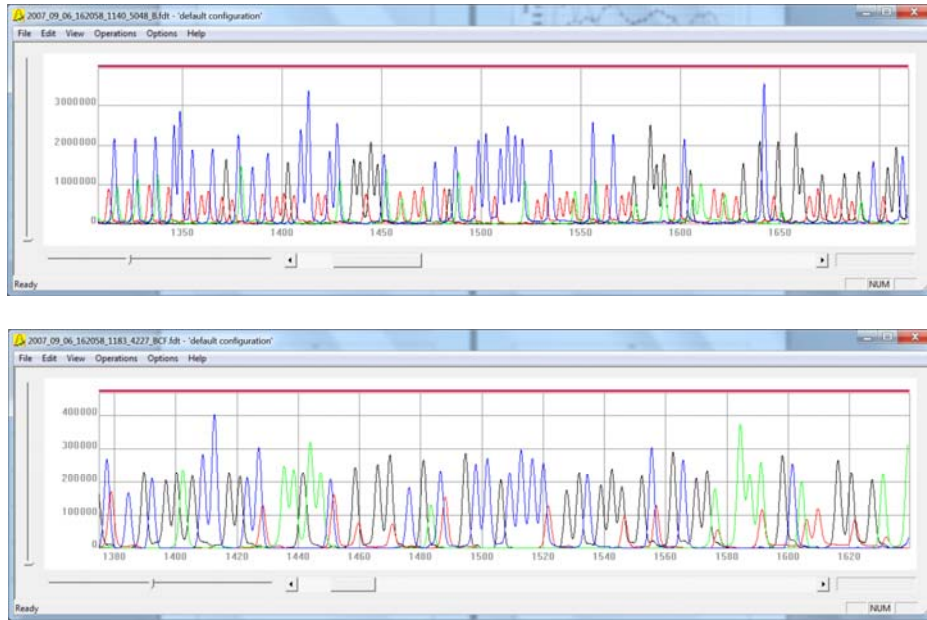


Figure 3.3.4. Sequencing traces processed using full spectrum analysis (upper panel) and virtual filter approach (lower panel).

As can be seen, the signal obtained from virtual filter approach is almost 10 fold lower. Another advantage of using the entire fluorescence spectrum obtained from the dye is a significant increase of the detection dynamic range. Indeed, in our 32-channel photosensor the

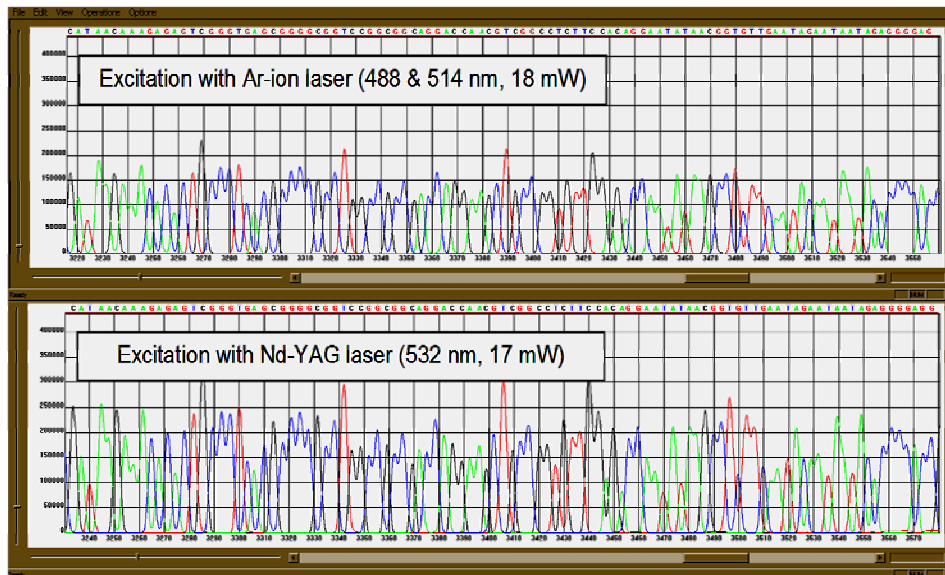


Figure 3.3.5. DNA sequencing traces obtained with Ar-ion and Nd-YAG lasers. channels have linear response up to $\sim 7 \times 10^7$ photocounts per second and if the measured

fluorescent dye has a broad spectrum, the linear range of the detection for this dye will be proportional to the number of simultaneously illuminated channels of the sensor.

Figure 3.3.5 presents two DNA sequencing traces obtained the developed data acquisition system using single photon spectrometer and two lasers.

We found that sequencing traces recorded with the two lasers look very similarly and the quality scores are nearly identical. The obtained sequence read length is as long as ~650 base pairs which is in a very good agreement with the read lengths obtained for the same capillary length in commercial DNA sequencers (ABI-3730, AppliedBiosystems Inc., Foster City, CA).

II. CONCLUSION

I have developed a universal data acquisition system for measurements of multi-color fluorescence in Life Science applications. The system fully satisfies the goal which has been set in the beginning of my work: Development of universal data acquisition system for ultra-sensitive and highly accurate detection of very weak fluorescent signals from biological objects.

Unique features of the developed system include:

- Supports wide variety of tests and chemistries
- Most detailed and accurate measurement of fluorescent signals enabling any desired data processing
- Single photon sensitivity
- Low reagent consumption (*low dead volume*)
- Ability of repetitive runs of the same samples in different conditions
- No equipment conservation needed
- Ease of use and maintenance
- Modular (plug & play) architecture

Application of the system in three most widely used areas have been demonstrated and tested:

- Multi-color fluorometry
- Detection and recognition of color-encoded beads
- DNA sequencing by capillary electrophoresis

We believe that due to an extremely high sensitivity of the photon detection in conjunction with method for highly accurate color decomposition, the developed system may be very useful for detection and decoding of multiple color codes and it can be used in various applications which require highly accurate identification of biological samples labeled with multiple fluorescent colors.

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